

Some relativistic elements

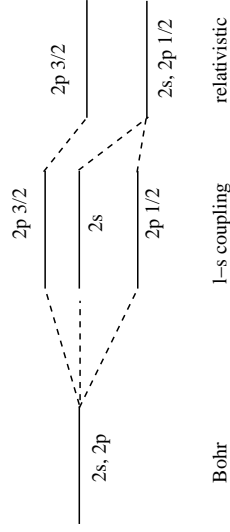
Peter Reinhardt

Laboratoire de Chimie Théorique, Université Paris VI, 75252 Paris CEDEX 05,
Peter.Reinhardt@upmc.fr

-p. 1

Motivation

- Why should we bother about relativity ? Organic chemistry deals with light atoms
- But already hydrogen: fine structure



- Heavy-element chemistry: nuclear waste, Pb-induced diseases, Pt catalysts
- Relativistic effects may be important already around $Z = 30$ (Zn)

-p. 2

-p. 3

-p. 4

Lorentz transformations

Two reference systems with relative speed v with respect to each other should maintain the same laws of physics, in particular the same speed of light c :

$$\vec{r}' = \frac{\vec{r} - \vec{v}t}{\sqrt{1 - v^2/c^2}}$$

$$t' = \frac{t - r v^2/c^2}{\sqrt{1 - v^2/c^2}}$$

$x^2 + y^2 + z^2 = c^2 t^2$ yields $(x')^2 + (y')^2 + (z')^2 = c^2 (t')^2$

- Limit of $v \rightarrow 0$ we find again $\vec{r}' = \vec{r} - \vec{v}t$ and $t' = t$ (transformation of Galilei)
- Consequence: muons arrive on earth despite their time of life of $2.2 \mu\text{s}$ (i.e. 660 m at speed c without $t \rightarrow t'$)
- Factor $\sqrt{1 - v^2/c^2}$ may be incorporated in the mass: $m = \frac{m_0}{\sqrt{1 - v^2/c^2}}$

-p. 5

-p. 6

Lorentz transformations

Relativistic energy expression

$$E^2 = m^2 c^4$$

$$= m^2 \left(1 - \frac{v^2}{c^2}\right) c^4 + m^2 v^2 c^2 = m_0^2 c^4 + p^2 c^2$$

de Broglie

- Particle without mass

$$E = pc = h\nu = h \frac{c}{\lambda} \longrightarrow p = \frac{h}{\lambda}$$

- Particle with mass

$$E = \sqrt{m_0^2 c^4 + p^2 c^2} = m_0 c^2 \sqrt{1 + \frac{p^2}{m_0^2 c^2}}$$

$$\approx m_0 c^2 \left(1 + \frac{1}{2} \frac{p^2}{m_0^2 c^2}\right) = E_M + \frac{p^2}{2m_0^2}$$

-p. 5

-p. 7

Bohr model

- Speed of an electron at level n

$$v_n = \frac{e^2}{4\pi\epsilon_0\hbar} \frac{Z}{n} = \alpha c \frac{Z}{n}$$

with the finestructure constant $\alpha = \frac{e^2 c}{4\pi\epsilon_0\hbar} = \frac{1}{137.037}$

- Orbital radius

$$r_n = \frac{\hbar^2}{m} \frac{4\pi\epsilon_0}{e^2} \frac{n^2}{Z} = \left(\sqrt{1 - \frac{v^2}{c^2}} \right) \frac{\hbar}{m_0} \frac{4\pi\epsilon_0}{e^2} \frac{n^2}{Z}$$

- Speed independent of the mass of the electron $\rightarrow v$ as a parameter

$$\sqrt{1 - \frac{v^2}{c^2}} = \sqrt{1 - \frac{\alpha^2 Z^2}{n^2}}$$

- Contraction of radii with Z and n .

-p. 8

-p. 9

Bohr model

element	Z	$1 - r'(1s)/r(1s)$	$1 - r'(4s)/r(4s)$
H	1	0.003 %	
Zn	30	2.5 %	0.15 %
Cd	48	6.5 %	0.38 %
Hg	80	19 %	1.07 %

- May be transferred to Slater's model. Internal contraction leaves more space for the external electrons \rightarrow effective n (4.7 for $n = 5$ etc.).
- Lanthanide contraction: higher Z and relativistic effects.
- s orbitals will be more touched than p or d , f
- s and p contracted, better screening, d expanded.

-p. 8

-p. 10

Relativistic quantum mechanics

- Schrödinger equation

$$\frac{p^2}{2m} = E \longrightarrow \frac{p^2}{2m}\psi = i\hbar\frac{\partial}{\partial t}\psi$$

- Derivatives in space and time in different degree \longrightarrow not Lorentz invariant
- Use $p^2c^2 + m^2c^4 = E^2$ instead:

$$\left(\nabla^2 - \frac{1}{c^2}\frac{\partial^2}{\partial t^2}\right)\phi = \frac{m^2c^2}{\hbar^2}\phi$$

(Klein-Gordon equation 1927, describes free particles with and without mass)

- But not an equation of first order in time and space!

-p. 11.

-p. 12.

Relativistic quantum mechanics

Take the square root of the $\nabla^2 - \frac{1}{c^2}\frac{\partial^2}{\partial t^2}$

$$\nabla^2 - \frac{1}{c^2}\frac{\partial^2}{\partial t^2} = (A\partial_x + B\partial_y + C\partial_z + \frac{i}{c}D\partial_t)(A\partial_x + B\partial_y + C\partial_z + \frac{i}{c}D\partial_t)$$

only possible with $AB + BA = 0, \dots$ and $A^2 = B^2 = \dots = 1$

$$(A\partial_x + B\partial_y + C\partial_z + \frac{i}{c}D\partial_t)\psi = \kappa\psi$$

or

$$\left(\nabla^2 - \frac{1}{c^2}\frac{\partial^2}{\partial t^2}\right)\psi = \kappa^2\psi$$

Identify $\kappa = \frac{mc}{\hbar} : (A\partial_x + B\partial_y + C\partial_z + \frac{i}{c}D\partial_t - \frac{mc}{\hbar})\psi = 0$

$(A, B, C) = i\beta\alpha_k$ and $D = \beta \longrightarrow$ Dirac equation

$$\left(\beta m_0 c^2 + \sum_{k=1}^3 \alpha_k p_k c\right) \psi(\vec{x}, t) = i\hbar \frac{\partial \psi(\vec{x}, t)}{\partial t}$$

-p. 11.

-p. 13.

Relativistic quantum mechanics

Pauli matrices

$$\sigma_k = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}, \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}, \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$$

lead to the 4×4 matrices

$$\alpha_k = \begin{pmatrix} 0 & \sigma_k \\ \sigma_k & 0 \end{pmatrix} \quad \beta = \begin{pmatrix} I_2 & 0 \\ 0 & -I_2 \end{pmatrix}$$

Dirac equation: (4-component matrix notation with two-component spinors ϕ_+ , ϕ_-):

$$\begin{pmatrix} mc^2 & c\sigma \cdot p \\ c\sigma \cdot p & -mc^2 \end{pmatrix} \begin{pmatrix} \phi_+ \\ \phi_- \end{pmatrix} = i\hbar \frac{\partial}{\partial t} \begin{pmatrix} \phi_+ \\ \phi_- \end{pmatrix}$$

Still only for one free particle, but relativistic

Relativistic quantum mechanics

- Dirac Hamiltonian with external fields

$$\hat{H}_D(i) = \left[q_i V(\vec{r}_i) + c \sum_{k=x,y,z} \alpha_k(i) (p_k - \frac{q}{c} A_k) + \beta m_0 c^2 \right]$$

V : electrostatic potential

\vec{A} : vector potential for magnetic fields $\vec{B} = \vec{\nabla} \times \vec{A}$

- Breit equation: add to the Dirac Hamiltonian an electron-electron coupling term

$$\left\{ \sum_i \hat{H}_D(i) + \sum_{i>j} \frac{1}{r_{ij}} - \sum_{i>j} \hat{B}_{ij} \right\} \Psi = E\Psi$$

with

$$\hat{B}_{ij} = -\frac{1}{2r_{ij}} \left[\vec{a}(i) \cdot \vec{a}(j) + \frac{(\vec{a}(i) \cdot \vec{r}_{ij})(\vec{a}(j) \cdot \vec{r}_{ij})}{r_{ij}^2} \right]$$

Relativistic quantum mechanics

Take only terms to order $(Z\alpha)^2$ (i.e. $Z < 136$)

$$H_{MV} = -\frac{\alpha^2}{8} \sum_i \Delta_i^2 \quad \text{Mass-velocity term}$$

$$H_{D1} = -\frac{\alpha^2 Z}{8} \sum_i \Delta_i \left(\frac{1}{r_i} \right) \quad \text{Darwin term}$$

$$H_{SO} = \frac{\alpha^2 Z}{2} \sum_i \frac{1}{r_i^3} \vec{l}_i \cdot \vec{s}_i \quad \text{Spin-orbit (of one electron)}$$

$$H_{D2} = -\frac{\alpha^2}{4} \sum_{i < j} \delta_i \left(\frac{1}{r_{ij}} \right) \quad \text{2-electron Darwin term}$$

$$H_{SOO} = -\frac{\alpha^2}{2} \sum_{i < j} \frac{\vec{r}_{ij} \times \vec{p}_i}{r_{ij}^3} (\vec{s}_i + 2\vec{s}_j) \quad \text{Spin-other-orbit}$$

Have to developed and applied to spinors \longrightarrow very heavy

-p.17

-p.19

In practice

The 4-component formalism may be reduced to a 1-component formalism for α (spin up) and β (spin down) spinors

- 1-electron terms:
 - Mass-velocity, Darwin term: spin independent
 - Spin-orbit: depends explicitly on the spin
- additional 1-electron integrals
- 2-electron terms: additional (or modified) 2-electron integrals

What can be incorporated in a pseudopotential ?

- A pseudopotential should not change the symmetry of an electronic state
- No way to have spin-orbit pseudopotentials
- But mass-velocity and Darwin terms can be included
- Treat spin-orbit coupling through CAS-SCF/MR-CI mixing all states which contribute to the final multiplets (atoms, diatomics)
- Closed-shell systems: no spin-orbit

-p.18

-p.20