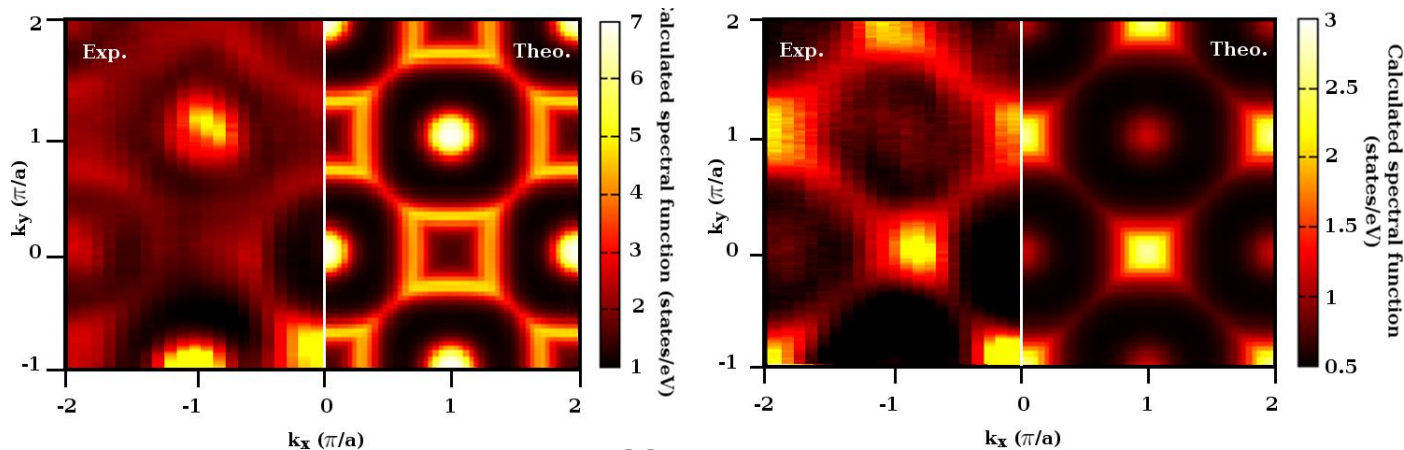


# The Quantum Many-Body Problem from a Dynamical Mean Field Theory Perspective

— or —

## Spectral Properties of Correlated Materials from First Principles



**Phys. Rev. Mat.**  
**2, 032001(R)**  
**(2018)**

**Silke Biermann**

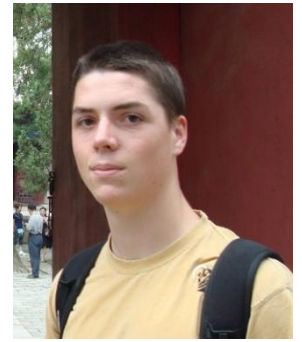
Centre de Physique Théorique,  
Ecole Polytechnique, Institut Polytechnique de Paris, France



**Cyril Martins**  
(now: MdC at  
Toulouse U.)

# “Electronic Structure Theory” at CPHT:

## Correlated materials from first principles



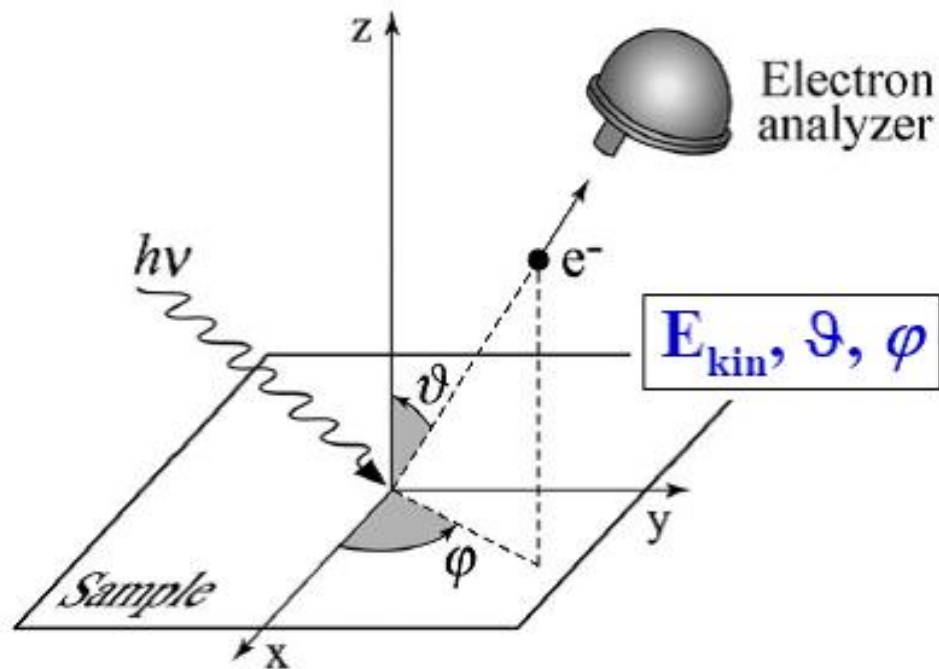
**A. van Roekeghem**  
(now: permanent  
researcher at CEA  
Grenoble)



S. Backes, J. Steinbauer, **Benjamin Lenz**, A. Galler, S. Bhandary, S. Panda, M. Turtulici  
(now MdC in Jussieu) Also: A. Subedi, L. Pourovskii

# Spectral properties ...

**Cf Pina's  
Talk !**

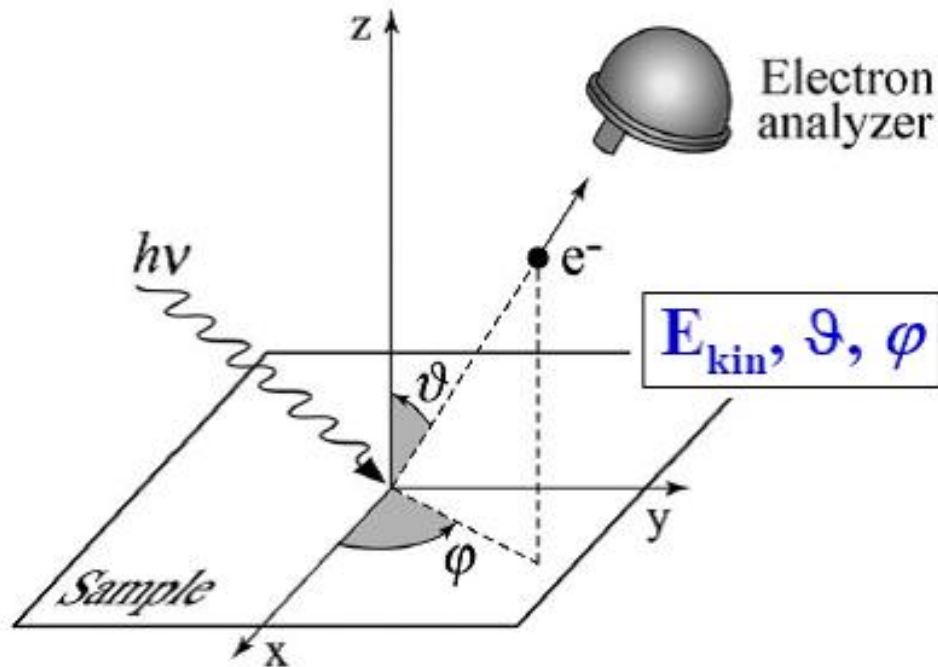


Can we calculate spectral functions from first principles ?

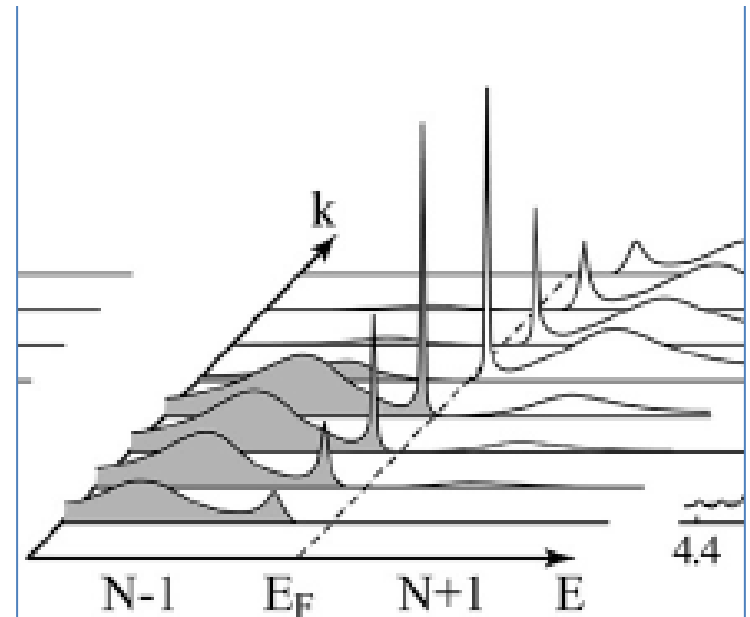
# Intuitive definition:

**Spectral function:  $A(k, \omega)$  describes electron addition and removal processes.**

**=> Photoemission and inverse photoemission**



Electron addition/removal spectra

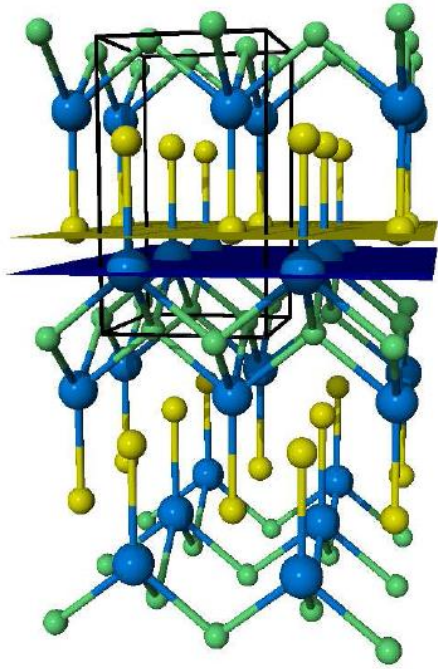


NB. Formal definition in language of Green's functions:  $A(k, \omega) = -1/\pi G(k, \omega)$

**Disclaimer: neglect matrix element effects.**



# Example: CeSF – an f-electron pigment



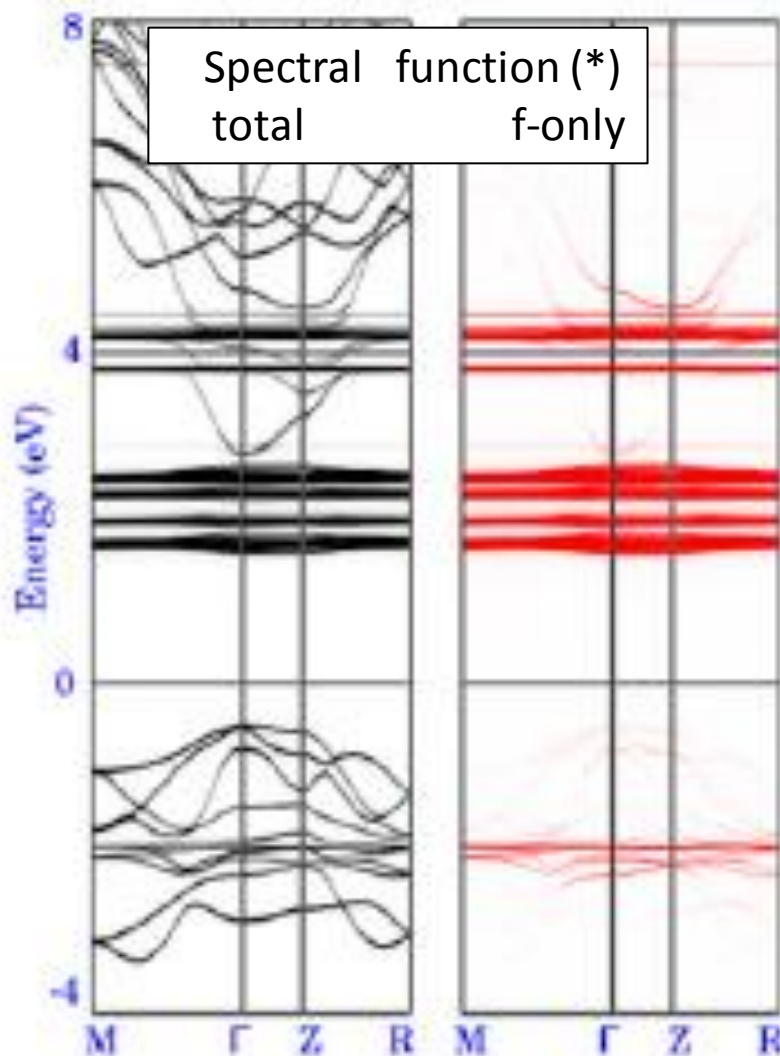
Ce:  $4f^1$  configuration,  
paramagnetic



Example :

CeSF – an f-electron pigment

(cf. Rhodia's Neolor series)



f-electrons in CeFS

- remain “atomic-like”, i.e. localized on the atomic sites
  - do not form bands
  - no Bloch states ...
- => “Mott insulator”
- Spectral feature: multiplets ...

Ce:  $4f^1$  configuration,  
paramagnetic

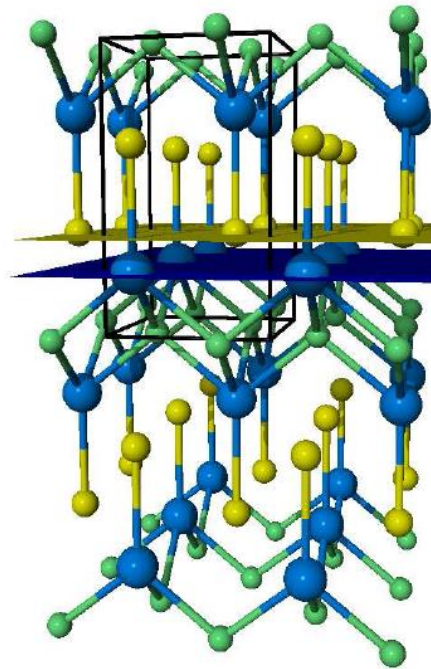
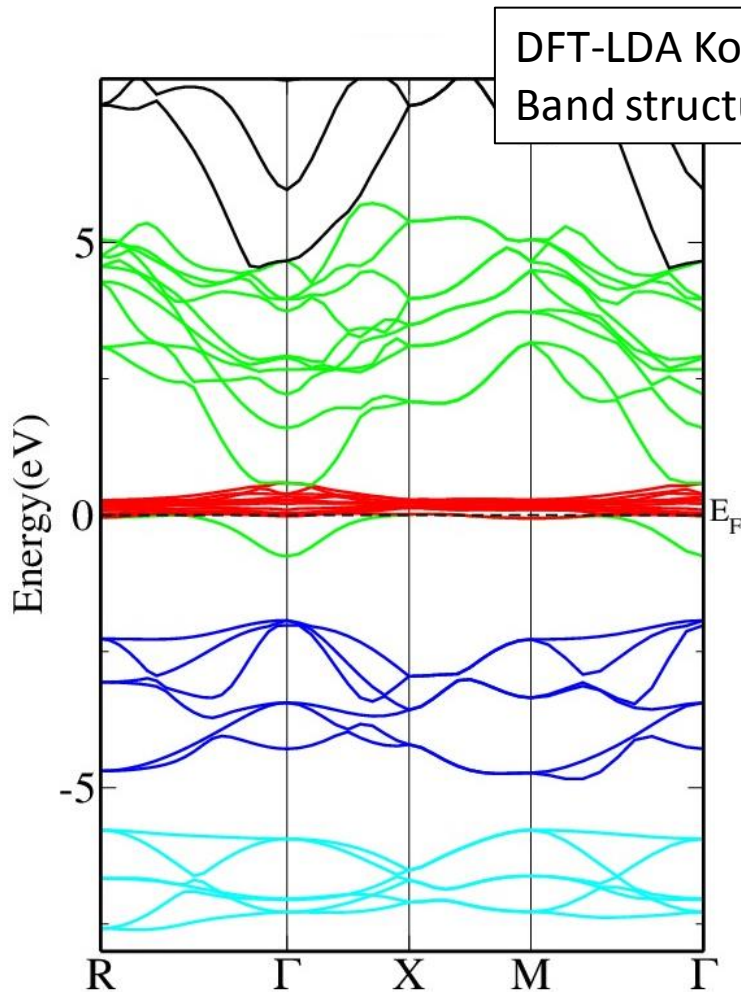
Calculated colour of CeSF:



Tomczak, Pourosvkii, Vaugier, Georges, Biermann, PNAS (2013)

(\*) calculated from Density Functional + Dynamical Mean Field Theory

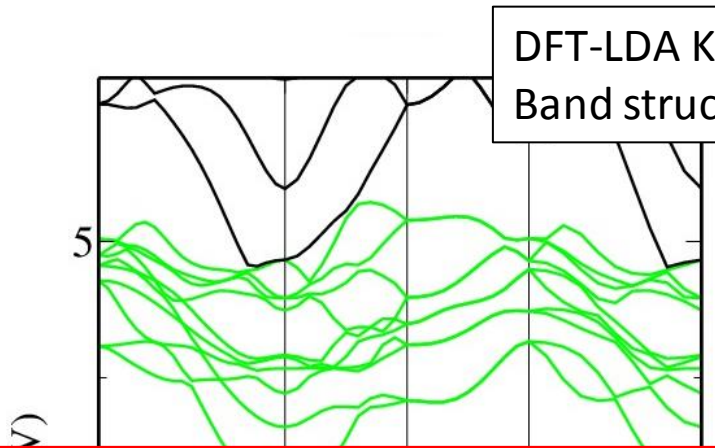
# Example: CeSF – an f-electron pigment



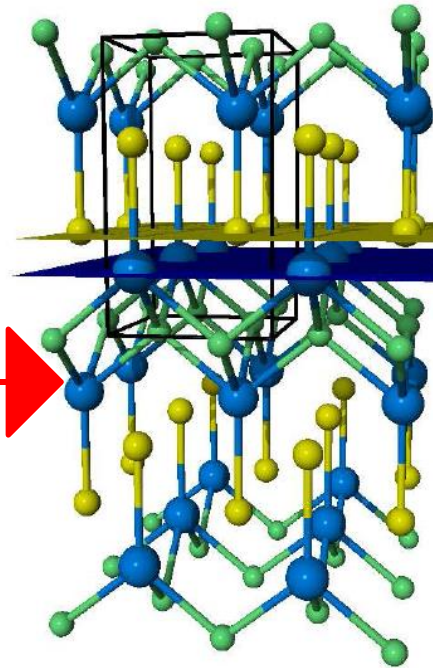
Ce:  $4f^1$  configuration,  
paramagnetic



# Example: CeSF – an f-electron pigment



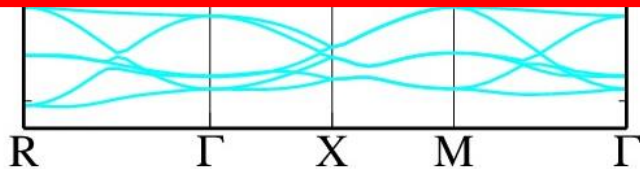
DFT-LDA Kohn-Sham  
Band structure



**Band structure not even  
qualitatively correct !**

**Need effective local Coulomb  
interaction (“Hubbard U”) on the Ce 4f  
shell**

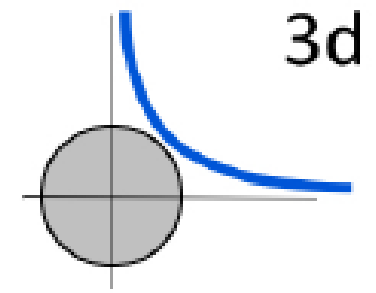
Ce:  $4f^1$  configuration,  
paramagnetic



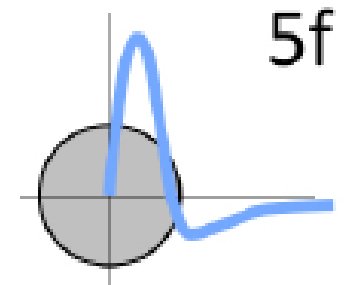


# Correlated materials ...

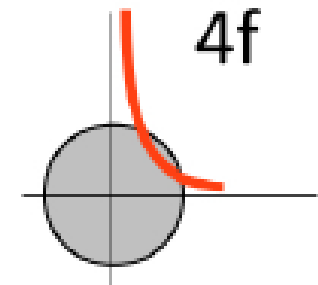
Examples: transition metal or rare-earth compounds (partially filled d- or f-shells), low-d systems ...



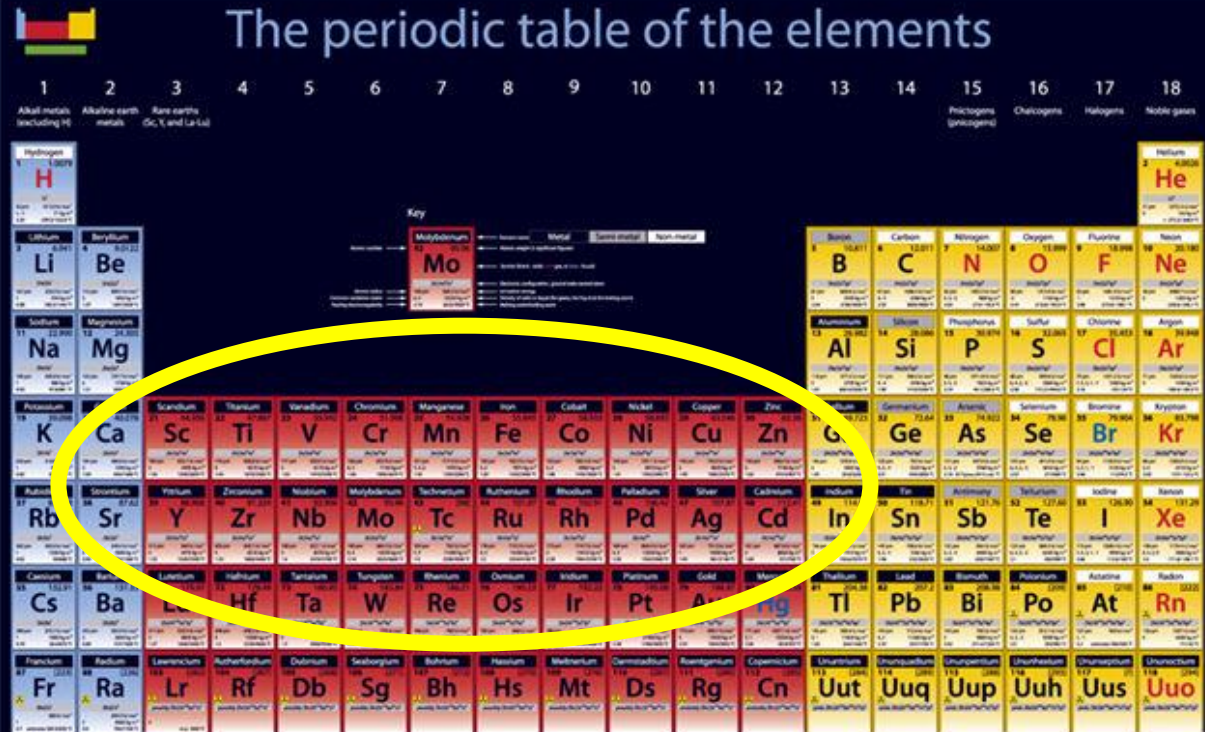
transition



actinide

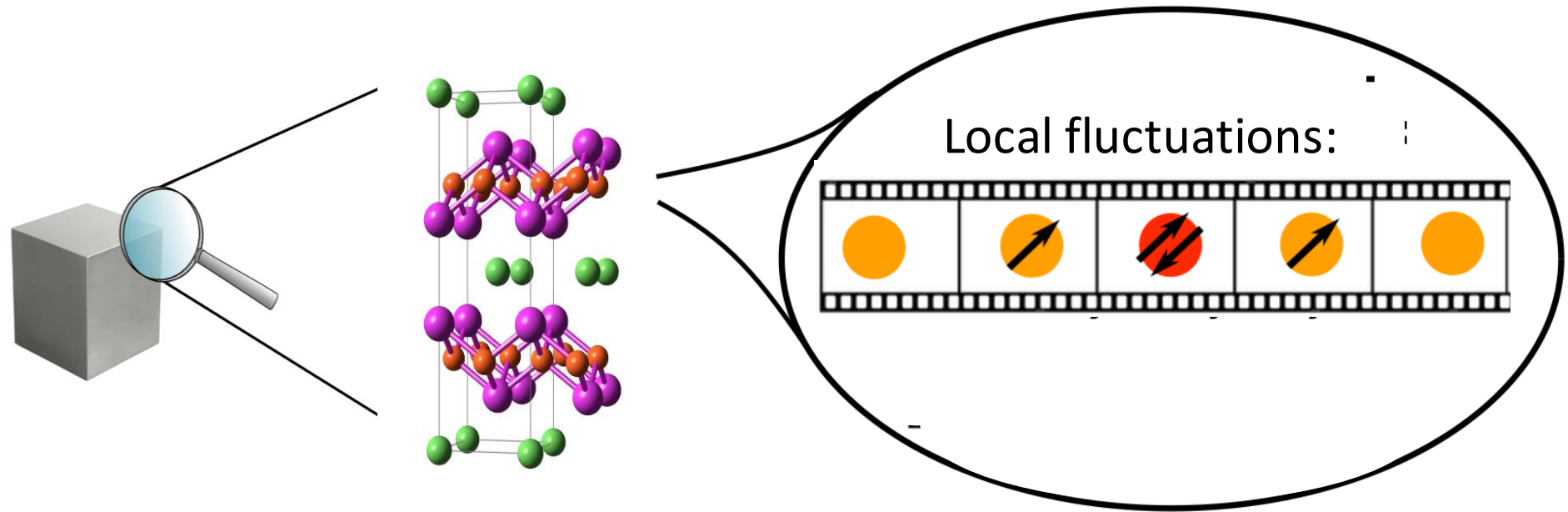


rare earth



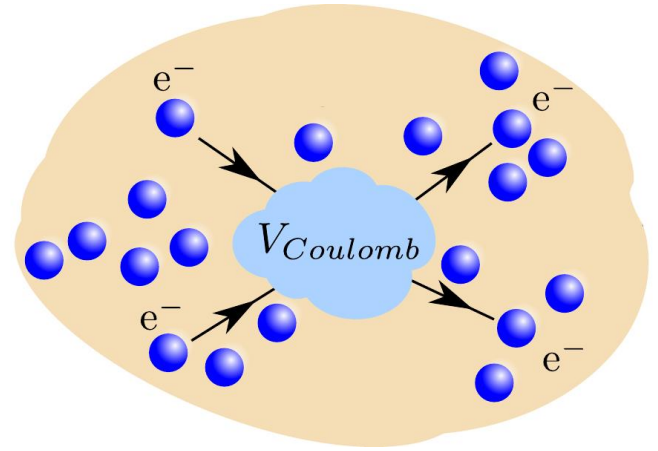
Radial wave function  $(n,l)$  has  $n-l-1$  nodes  
=> Spatial extension of 3d and 4f small.

Most challenging regime is in between  
the two (band and atomic) limits !



# Correlated electron materials:

**Coulomb interactions invalidate a simple single-particle picture**



Correlated materials

... display **collective behavior** of the electrons (beyond single-particle picture), **emergent properties**

... are characterized by various **competing energy scales, competing instabilities**

... are **extremely sensitive** with respect to external parameters

... need theoretical treatment beyond band theory

# First Principles Calculations for Correlated Materials

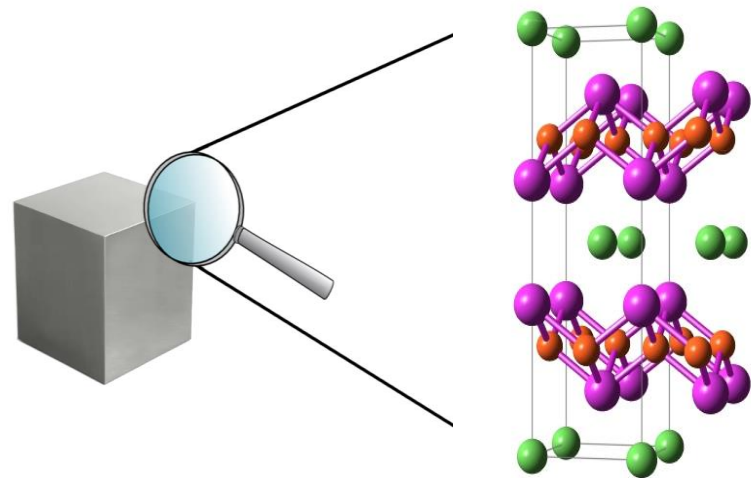
Calculate properties of **materials with strong electronic Coulomb correlations**

- ground- and excited state properties (spectra, optics, correlation functions ...)
  - beyond the single-particle picture (“beyond mean field”)
  - finite temperatures
  - from *first principles*, i.e. without adjustable parameters
- ⇒ Strategy: combine techniques from **many-body theory** & **first principles electronic structure theories**

Here: **Dynamical Mean Field Theory**  
**(DMFT)**



# Dynamical Mean Field Theory

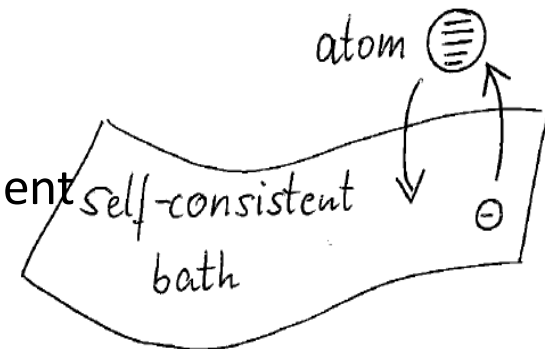


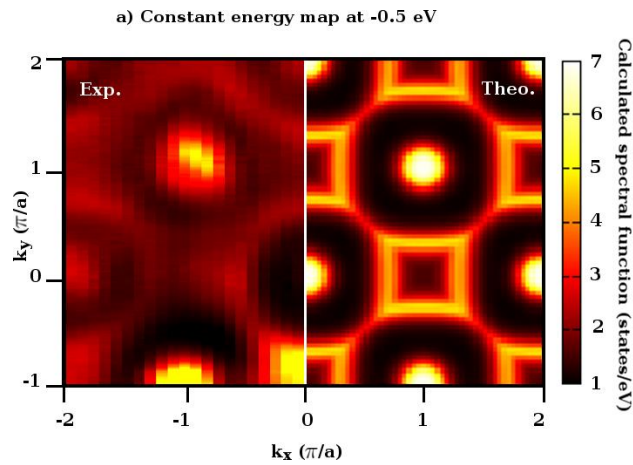
**Dynamical mean-field theory (DMFT):**  
Includes local dynamical quantum fluctuations in a many-body way



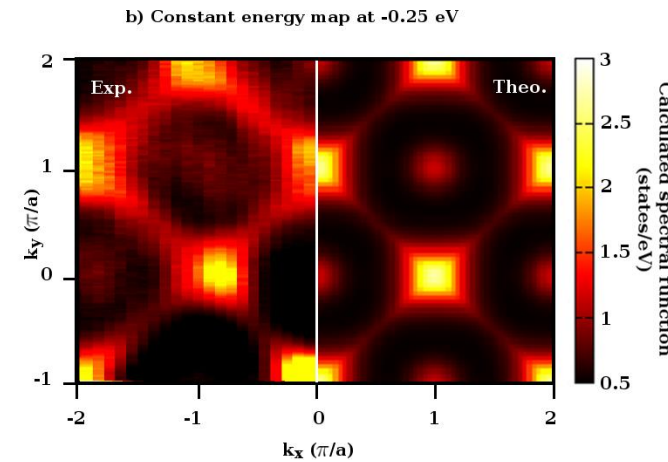
A. Georges, et al., Rev. Mod. Phys. 68, 13 (1996)

More technically speaking:  
DMFT approximates dynamical orbital- and momentum-dependent many-body self-energy by a momentum-independent self-energy, calculated from a self-consistent effective local quantum impurity problem





## 2<sup>nd</sup> example: Sr<sub>2</sub>IrO<sub>4</sub>

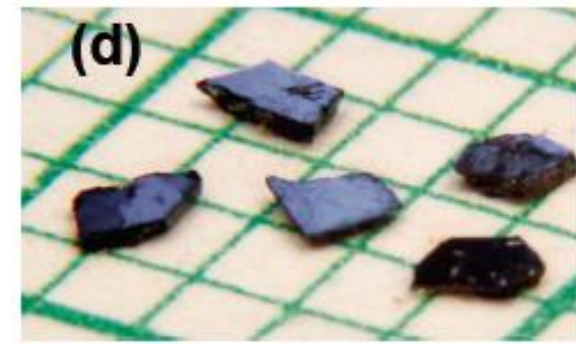
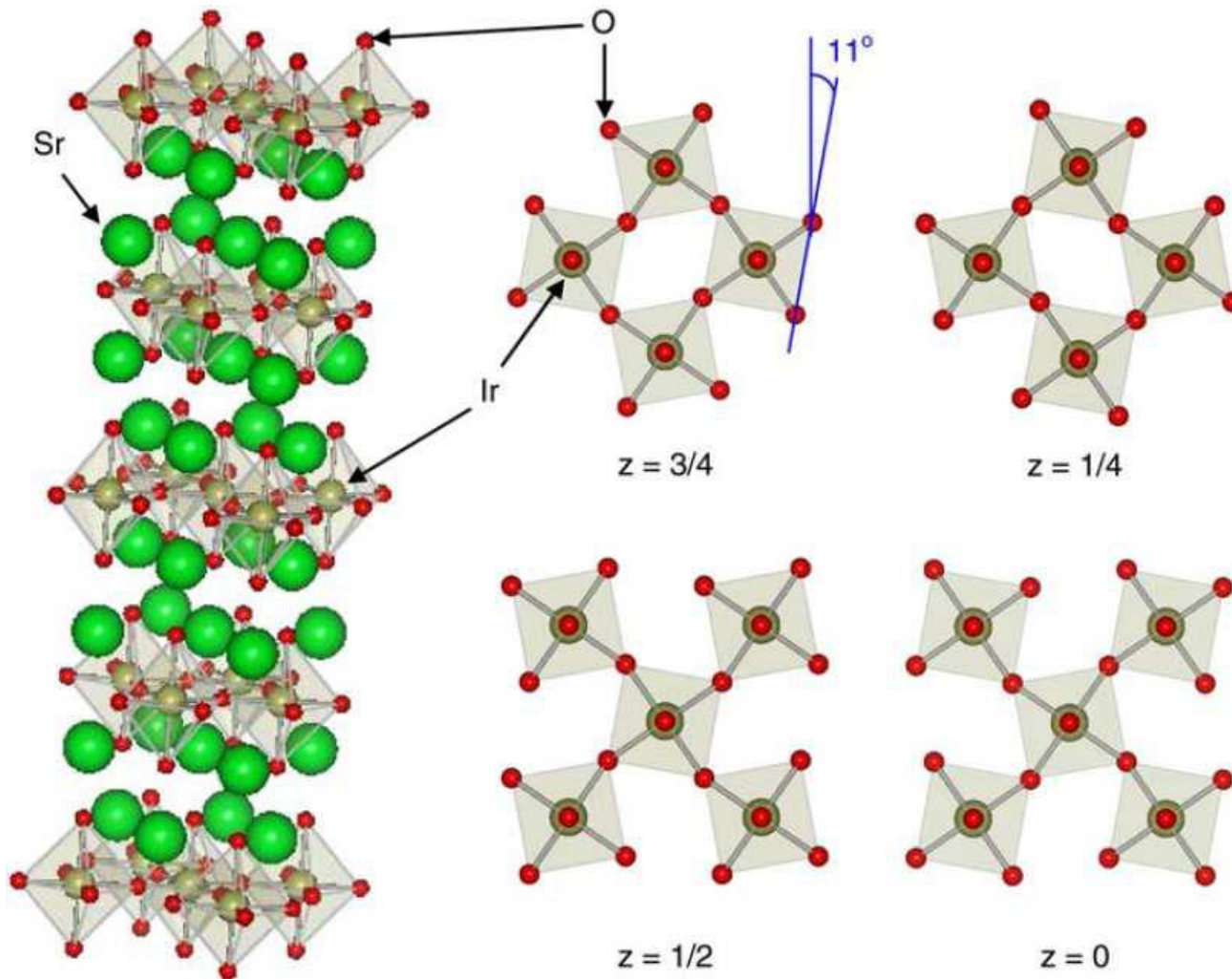


- Combined Density Functional Theory + Dynamical Mean Field Theory for Sr<sub>2</sub>IrO<sub>4</sub> and Sr<sub>2</sub>RhO<sub>4</sub>
- Non-local correlations in Sr<sub>2</sub>IrO<sub>4</sub>
- Spectral properties of doped Sr<sub>2</sub>IrO<sub>4</sub> – a simple picture for the pseudo-gap

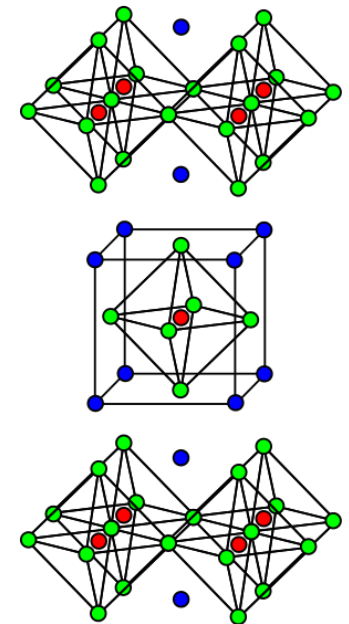
# Recent literature (not exhaustive) ....

- Kim et al, Cao et al.,
- Moon et al
- Korneta et al.,
- Nichols et al.
- Li et al.,
- Wang et al,
- Lu et al.
- Moretti et al.,
- Lovesey et al.
- Yamasaki et al.
- Boseggia et al.
- Calder et al.
- Ronnow et al.
- Fujiyama et al.
- Zocco et al.
- Jackeli and Khaliullin
- Carter and Kee
- Kim, Daghofer et al.
- Paerschke et al.
- Arita et al.,
- Watanabe et al.
- Zhang et al
- Scheiderer et al.
- Hampel et al.,
- Proepper et al.
- Agrestini et al.
- Cheng et al.
- Chikara et al.
- Chatterje et al.
- Moutenet et al.
- Chen et al.
- Zhao et al.
- Cao et al.
- Lu et al.
- Wang et al.
- Guevara et al.
- .....

# Sr<sub>2</sub>IrO<sub>4</sub>: K<sub>2</sub>NiF<sub>4</sub>-structure (cf. La<sub>2</sub>CuO<sub>4</sub>, Sr<sub>2</sub>(Rh/Ru)O<sub>4</sub>)



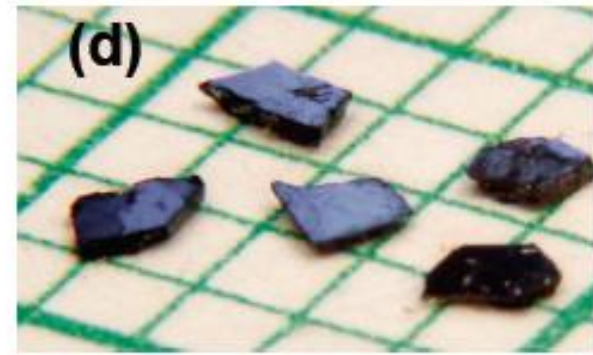
(Qi et al, PRB '12)



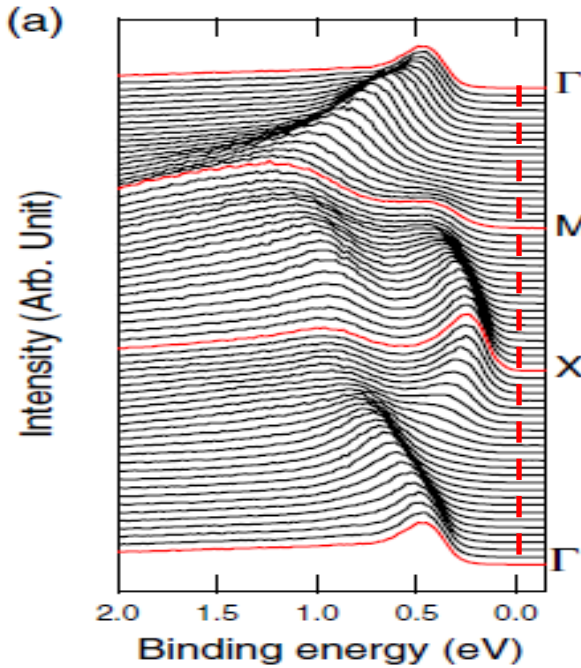
Ideal double perovskite structure



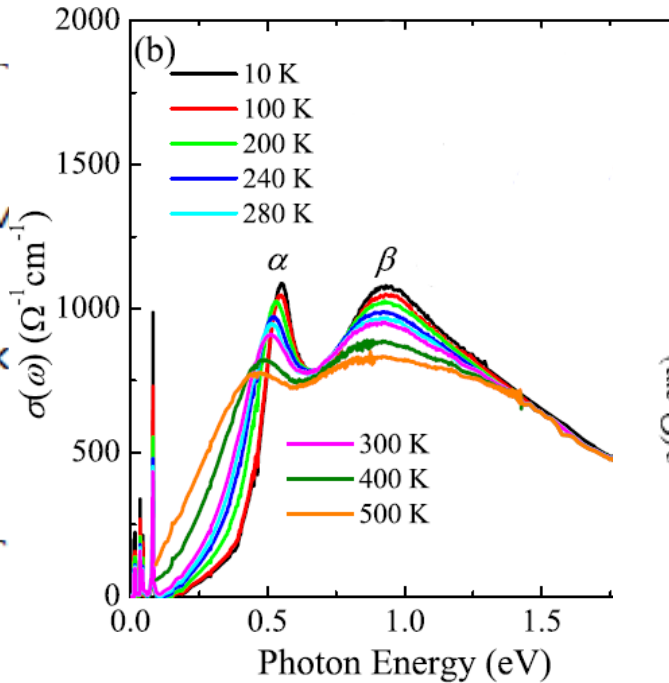
# Sr<sub>2</sub>IrO<sub>4</sub>: ARPES, Optics, Transport ...



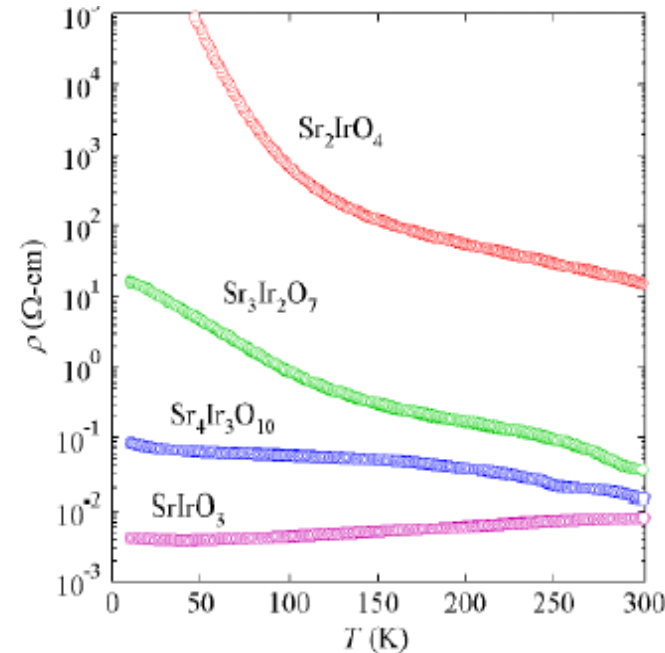
Qi et al, Phys. Rev. B 86, 125105 (2012)



Kim et al, Phys Rev Lett 101, 076402 (2009)



Moon et al, Phys Rev B 80, 195110 (2009)

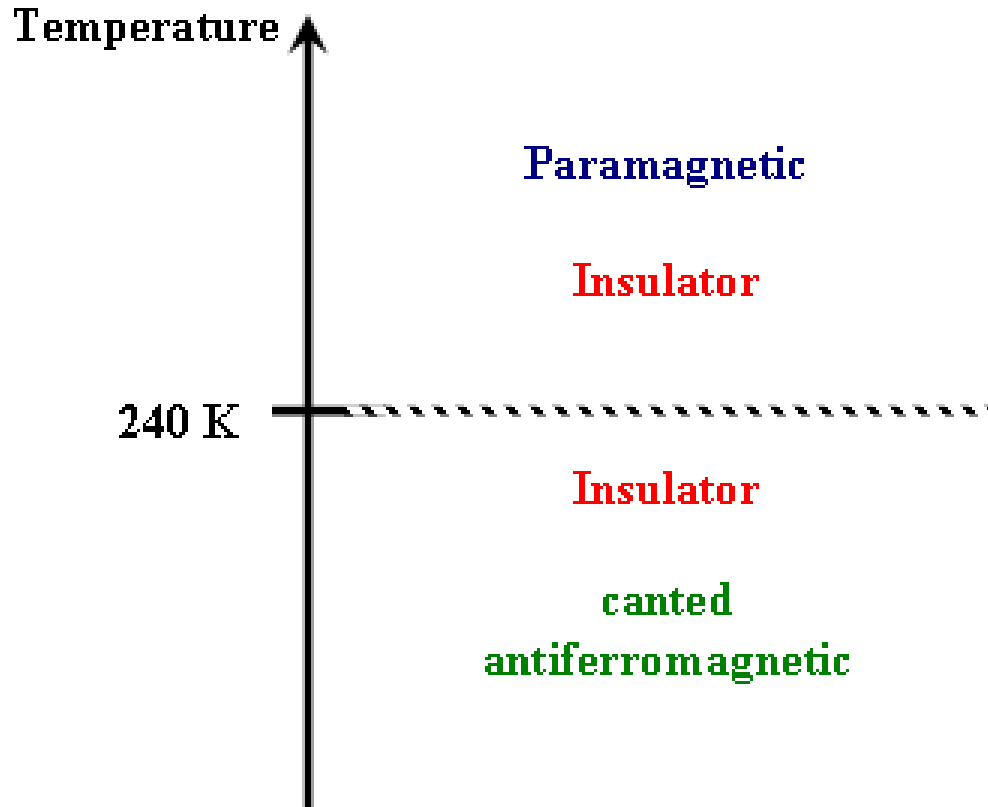


J. Matsuno H. Takagi et al, PRL (2015)

⇒ insulating behavior at all measured temperatures (up to ~ 500K).  
 Strongly T-dependent gap.

⇒ Electronic configuration: 5d<sup>5</sup> ⇒ not a band insulator.

# Sr<sub>2</sub>IrO<sub>4</sub>: Magnetism



← Our focus !

Below 240K:

Canted AFM

Weak ferromagnetic  
moment in ab-plane

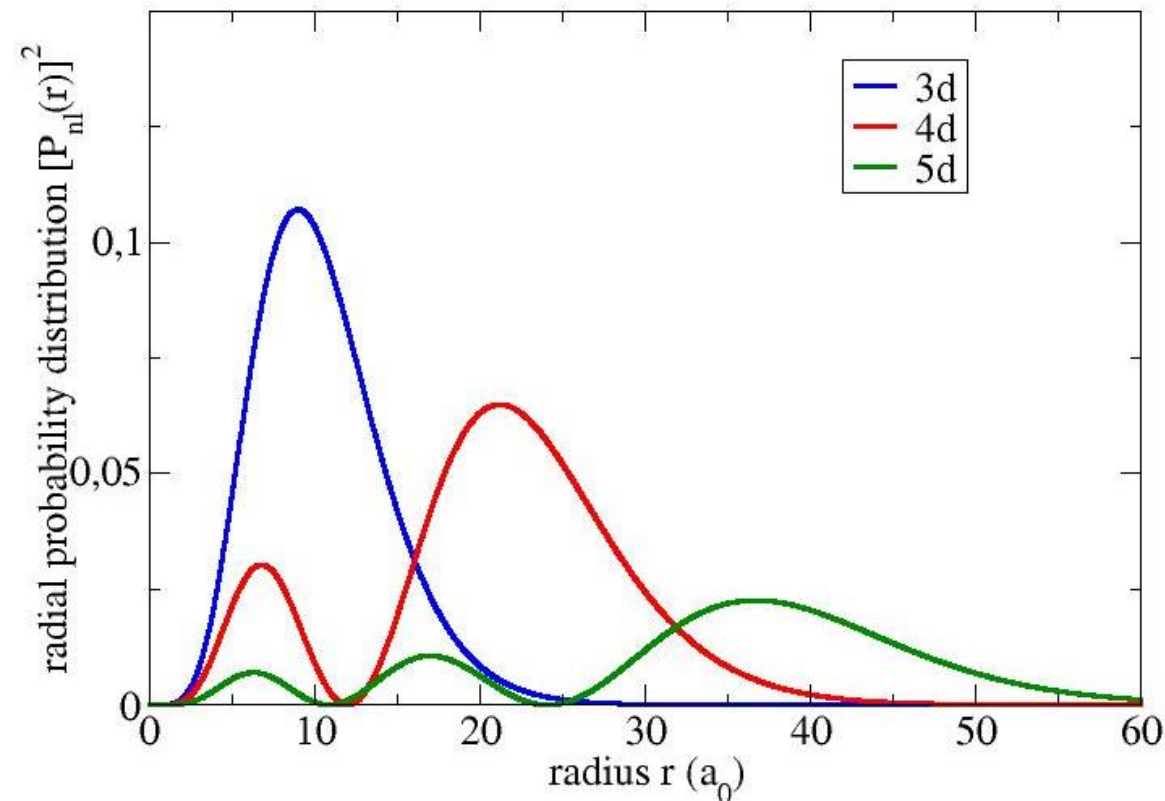
# From 3d to 4d to 5d ...

21 <b>Sc</b> [Ar]4s <sup>1</sup> 3d <sup>2</sup> scandium 44.96	22 <b>Ti</b> [Ar]4s <sup>1</sup> 3d <sup>3</sup> titanium 47.88	23 <b>V</b> [Ar]4s <sup>1</sup> 3d <sup>4</sup> vanadium 50.94	24 <b>Cr</b> [Ar]4s <sup>1</sup> 3d <sup>5</sup> chromium 52.00	25 <b>Mn</b> [Ar]4s <sup>2</sup> 3d <sup>5</sup> manganese 54.94	26 <b>Fe</b> [Ar]4s <sup>2</sup> 3d <sup>6</sup> iron 55.85	27 <b>Co</b> [Ar]4s <sup>2</sup> 3d <sup>7</sup> cobalt 58.93	28 <b>Ni</b> [Ar]4s <sup>2</sup> 3d <sup>8</sup> nickel 58.69	29 <b>Cu</b> [Ar]4s <sup>1</sup> 3d <sup>10</sup> copper 63.55	30 <b>Zn</b> [Ar]4s <sup>2</sup> 3d <sup>10</sup> zinc 65.39
39 <b>Y</b> [Kr]5s <sup>2</sup> 4d <sup>1</sup> yttrium 88.91	40 <b>Zr</b> [Kr]5s <sup>1</sup> 4d <sup>3</sup> zirconium 91.22	41 <b>Nb</b> [Kr]5s <sup>1</sup> 4d <sup>4</sup> niobium 92.91	42 <b>Mo</b> [Kr]5s <sup>1</sup> 4d <sup>5</sup> molybdenum 95.94	43 <b>Tc</b> [Kr]5s <sup>2</sup> 4d <sup>5</sup> technetium (98)	44 <b>Ru</b> [Kr]5s <sup>2</sup> 4d <sup>6</sup> ruthenium 101.1	45 <b>Rh</b> [Kr]5s <sup>2</sup> 4d <sup>7</sup> rhodium 102.9	46 <b>Pd</b> [Kr]4d <sup>10</sup> palladium 106.4	47 <b>Ag</b> [Kr]5s <sup>1</sup> 4d <sup>10</sup> silver 107.9	48 <b>Cd</b> [Kr]5s <sup>2</sup> 4d <sup>10</sup> cadmium 112.4
57 <b>La*</b> [Xe]6s <sup>2</sup> 5d <sup>1</sup> lanthanum 138.9	72 <b>Hf</b> [Xe]6s <sup>2</sup> 4f <sup>14</sup> 5d <sup>2</sup> hafnium 178.5	73 <b>Ta</b> [Xe]6s <sup>2</sup> 4f <sup>14</sup> 5d <sup>3</sup> tantalum 180.9	74 <b>W</b> [Xe]6s <sup>2</sup> 4f <sup>14</sup> 5d <sup>4</sup> tungsten 183.9	75 <b>Re</b> [Xe]6s <sup>2</sup> 4f <sup>14</sup> 5d <sup>5</sup> rhenium 186.2	76 <b>Os</b> [Xe]6s <sup>2</sup> 4f <sup>14</sup> 5d <sup>6</sup> osmium 190.2	77 <b>Ir</b> [Xe]6s <sup>2</sup> 4f <sup>14</sup> 5d <sup>7</sup> iridium 190.2	78 <b>Pt</b> [Xe]6s <sup>1</sup> 4f <sup>14</sup> 5d <sup>9</sup> platinum 195.1	79 <b>Au</b> [Xe]6s <sup>1</sup> 4f <sup>14</sup> 5d <sup>10</sup> gold 197.0	80 <b>Hg</b> [Xe]6s <sup>2</sup> 4f <sup>14</sup> 5d <sup>10</sup> mercury 200.5

3d

4d

5d



Orbital extension  
increases

⇒ Weaker  
correlations

⇒ Unless .....

# From 3d to 4d to 5d ...

21 Sc [Ar]4s <sup>1</sup> 3d <sup>2</sup> scandium 44.96	22 Ti [Ar]4s <sup>1</sup> 3d <sup>3</sup> titanium 47.88	23 V [Ar]4s <sup>1</sup> 3d <sup>4</sup> vanadium 50.94	24 Cr [Ar]4s <sup>1</sup> 3d <sup>5</sup> chromium 52.00	25 Mn [Ar]4s <sup>2</sup> 3d <sup>5</sup> manganese 54.94	26 Fe [Ar]4s <sup>2</sup> 3d <sup>6</sup> iron 55.85	27 Co [Ar]4s <sup>2</sup> 3d <sup>7</sup> cobalt 58.93	28 Ni [Ar]4s <sup>2</sup> 3d <sup>8</sup> nickel 58.69	29 Cu [Ar]4s <sup>1</sup> 3d <sup>10</sup> copper 63.55	30 Zn [Ar]4s <sup>2</sup> 3d <sup>10</sup> zinc 65.39
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3d

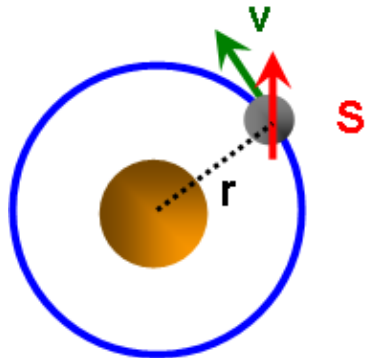
4d

5d

Cobalt (Z= 27)	$\zeta_{SO} \sim 0.06$ eV
Rhodium (Z= 45)	$\zeta_{SO} \sim 0.19$ eV
Iridium (Z= 77)	$\zeta_{SO} \sim 0.4$ eV

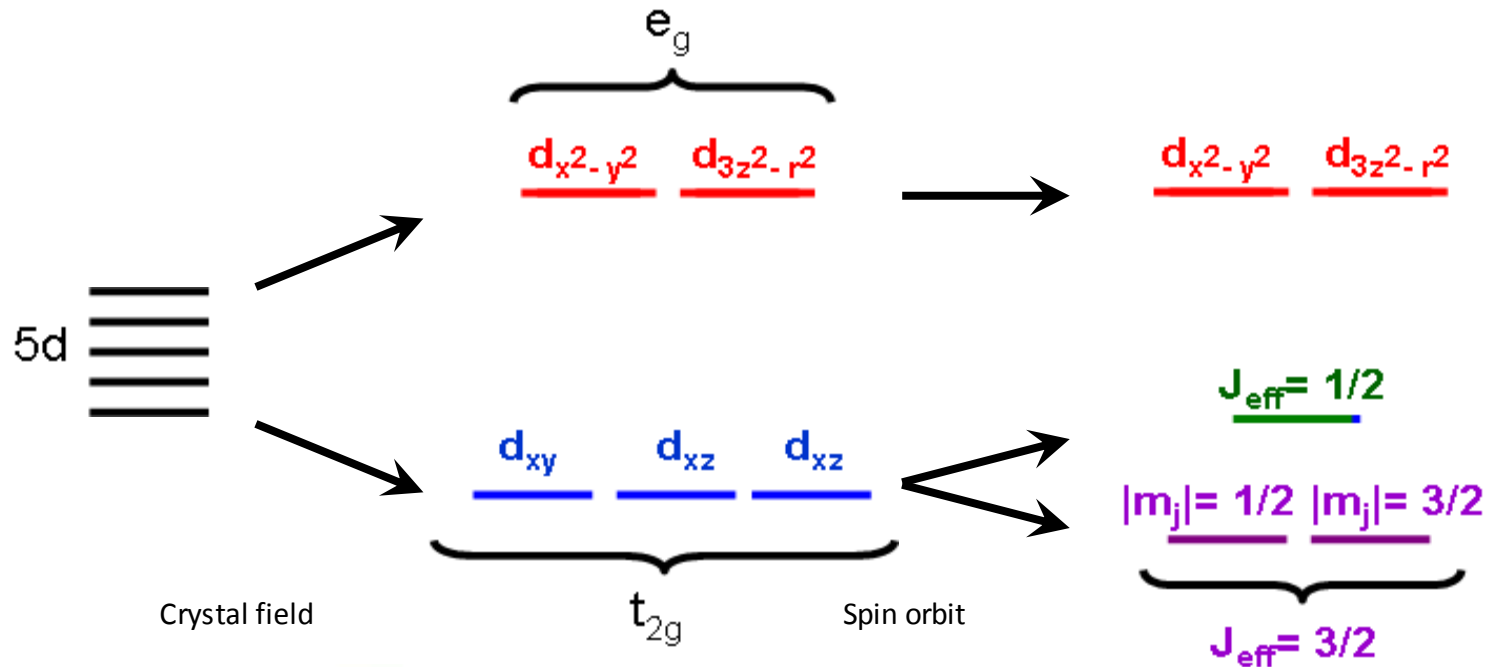
Atomic number increases

⇒ stronger spin-orbit coupling

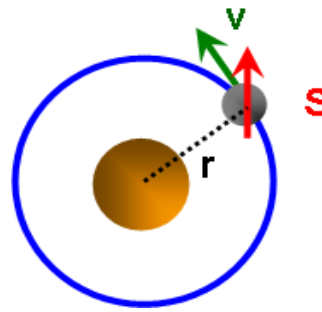
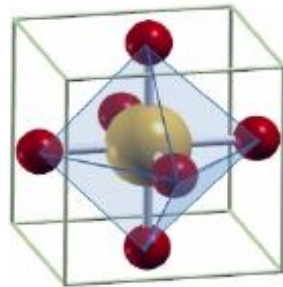




# Crystal field and spin-orbit: Ir in octahedral symmetry

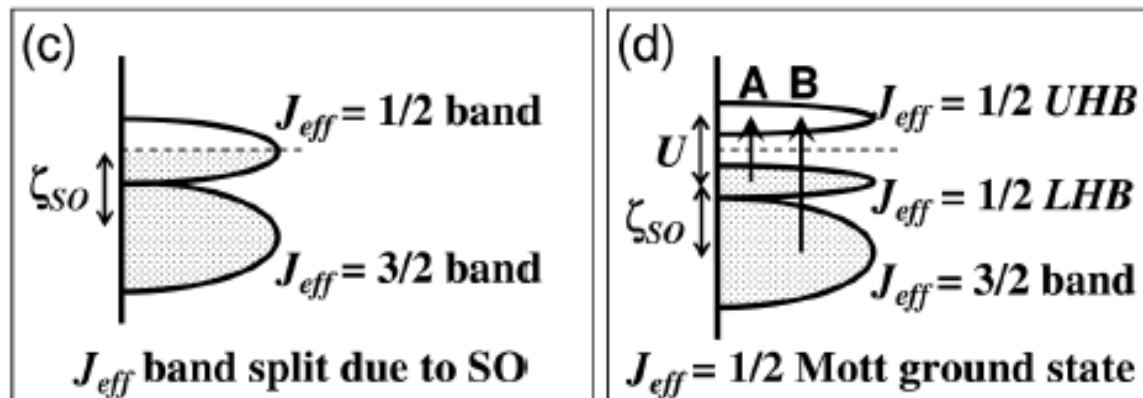


Crystal field



## Novel $J_{\text{eff}} = 1/2$ Mott State Induced by Relativistic Spin-Orbit Coupling in $\text{Sr}_2\text{IrO}_4$

B. J. Kim,<sup>1</sup> Hosub Jin,<sup>1</sup> S. J. Moon,<sup>2</sup> J.-Y. Kim,<sup>3</sup> B.-G. Park,<sup>4</sup> C. S. Leem,<sup>5</sup> Jaejun Yu,<sup>1</sup> T. W. Noh,<sup>2</sup> C. Kim,<sup>5</sup> S.-J. Oh,<sup>1</sup>  
J.-H. Park,<sup>3,4,\*</sup> V. Durairaj,<sup>6</sup> G. Cao,<sup>6</sup> and E. Rotenberg<sup>7</sup>



# Phase-Sensitive Observation of a Spin-Orbital Mott State in $\text{Sr}_2\text{IrO}_4$

B. J. Kim,<sup>1,2\*</sup> H. Ohsumi,<sup>3</sup> T. Komesu,<sup>3</sup> S. Sakai,<sup>3,4</sup> T. Morita,<sup>3,5</sup> H. Takagi,<sup>1,2\*</sup> T. Arima<sup>3,6</sup>

RIXS => " $J_{\text{eff}}=1/2$ " state

# DFT-LDA



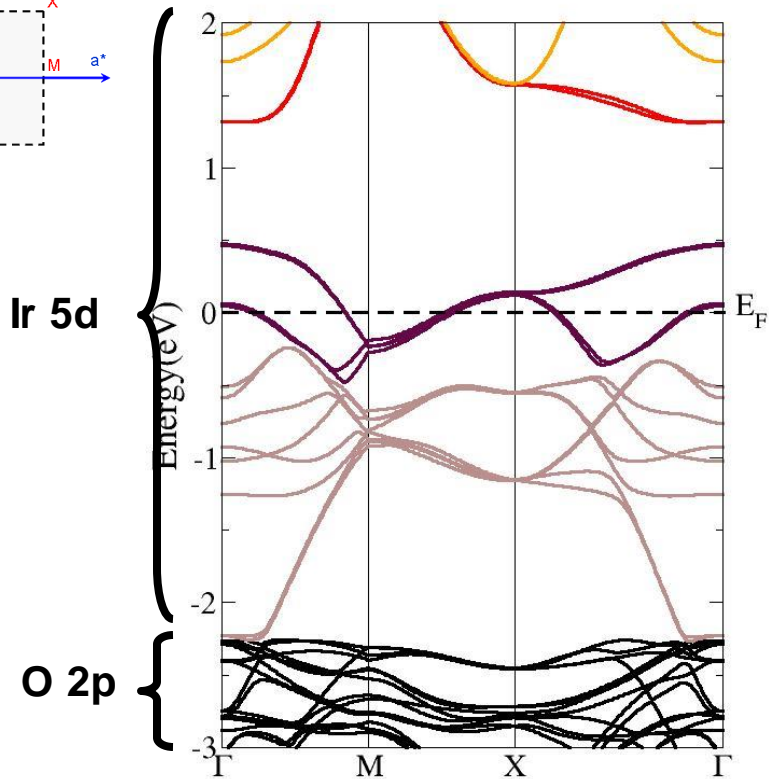
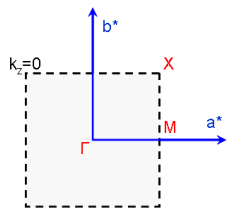
DFT part

Solving  
Kohn-Sham equations



$$\left[ -\frac{\hbar^2}{2m_0} \nabla^2 + V_{KS}(\mathbf{r}) \right] \psi_{\mathbf{k}}(\mathbf{r}) = \varepsilon_{n\mathbf{k}} \psi_{\mathbf{k}}(\mathbf{r})$$

within the Local Density Approximation (LDA)



}  $e_g$

4 bands cross the Fermi level:  
a **metallic** Kohn-Sham band structure

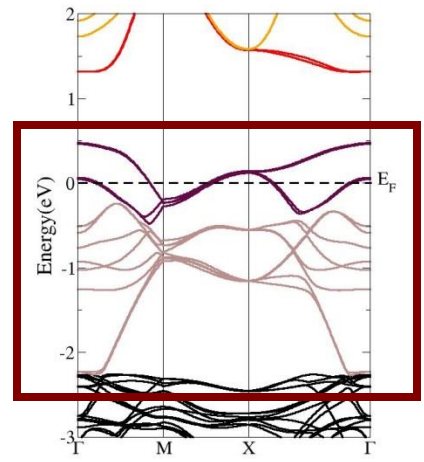
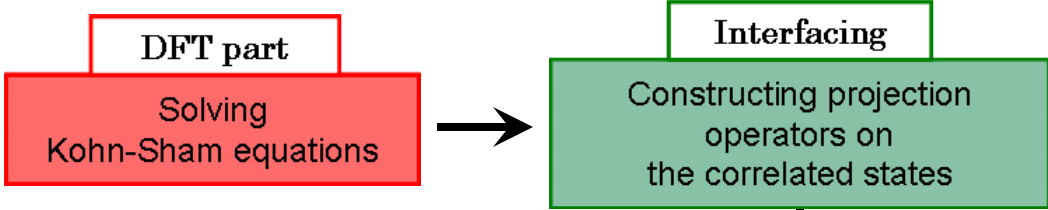
}  $j_{\text{eff}} = 1/2$   
 $j_{\text{eff}} = 3/2$

Used as a starting point for **LDA+DMFT** calculations.

*Martins et al, Phys. Rev. Lett. 107, 266404 (2011)*  
*(see also: Arita et al, Phys. Rev. Lett. 108, 086403 (2012),*  
*Zhang et al., Phys Rev Lett. 2013,*  
*Moutenet et al, Phys Rev B 2018)*

# Projection on the the $j_{\text{eff}} = 1/2$ and $3/2$ orbitals

Martins et al., Phys. Rev. Lett. 107, 266404 (2011)  
 J. Phys. Cond. Matt. (2017).

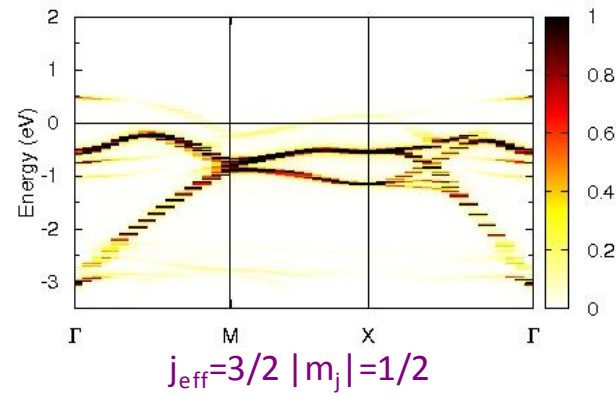
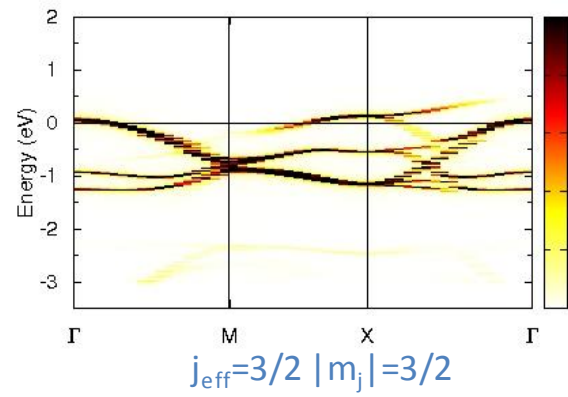
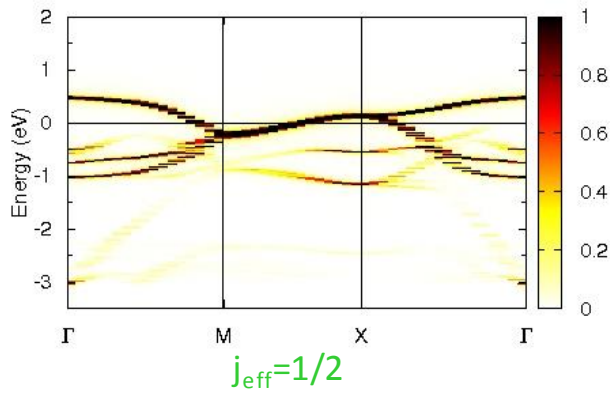


$$\hat{P}^\alpha(\mathbf{k}) = \sum_{m \in \mathcal{C}} |w_{\mathbf{k}m}^\alpha\rangle \langle w_{\mathbf{k}m}^\alpha|$$

Wannier-like orbitals obtained with a projection scheme

*Aichhorn et al, Phys Rev B 80, 085101 (2009)*

The  $j_{\text{eff}}=1/2$  state is almost **half-filled** with a bandwidth of  $\sim 1.5$  eV  
 the  $j_{\text{eff}}=3/2$  states are almost completely filled.



# Reduction of Effective Degeneracy by correlations

At the level of the DFT band structure, Sr<sub>2</sub>IrO<sub>4</sub> is not (yet) a single-orbital system:

---

	Sr <sub>2</sub> IrO <sub>4</sub>		
Wannier orbitals	$\left  \frac{1}{2}, \pm \frac{1}{2} \right\rangle$	$\left  \frac{3}{2}, \pm \frac{1}{2} \right\rangle$	$\left  \frac{3}{2}, \pm \frac{3}{2} \right\rangle$
occupation (LDA + SO)	1.16	1.98	1.84
charge (LDA + SO+DMFT)	1.02	2.00	1.98

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PRL 107, 266404 (2011)

PHYSICAL REVIEW LETTERS

week ending  
23 DECEMBER 2011

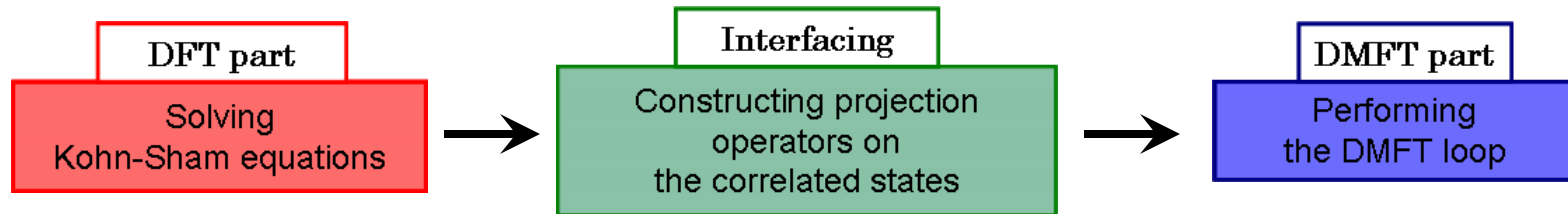
## Reduced Effective Spin-Orbital Degeneracy and Spin-Orbital Ordering in Paramagnetic Transition-Metal Oxides: Sr<sub>2</sub>IrO<sub>4</sub> versus Sr<sub>2</sub>RhO<sub>4</sub>

Cyril Martins,<sup>1,2</sup> Markus Aichhorn,<sup>3,1</sup> Loïg Vaugier,<sup>1</sup> and Silke Biermann<sup>1,2</sup>



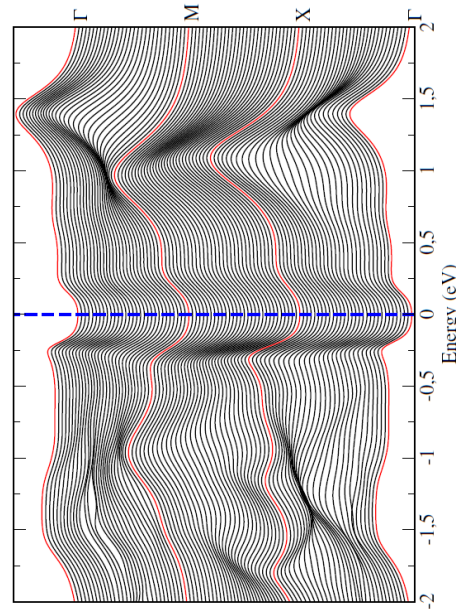
# Sr<sub>2</sub>IrO<sub>4</sub> a spin-orbital ordered Mott insulator

Martins et al., Phys. Rev. Lett. 107, 266404 (2011)



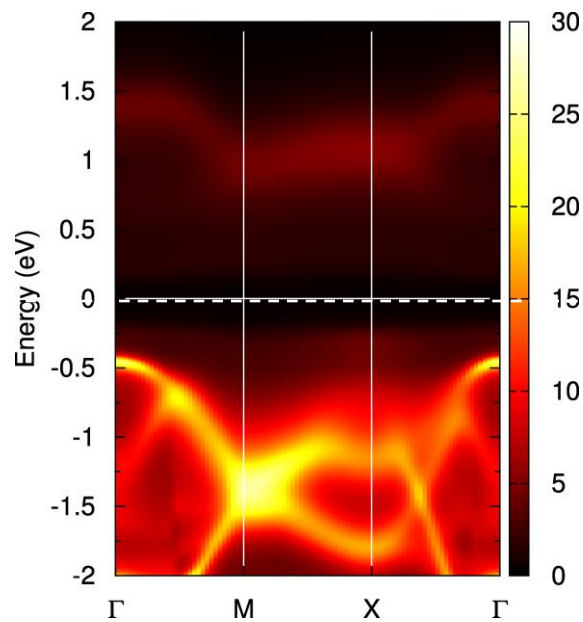
An **insulating paramagnetic state** with a gap of about **0.2 – 0.3 eV** is obtained. ( $U = 2.2$  eV and  $J = 0.3$  eV)

Jeff=1/2 spectral function



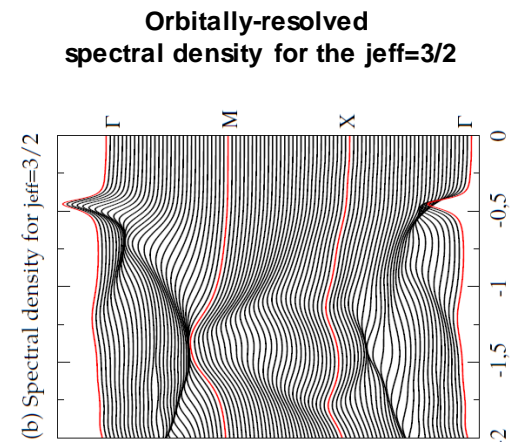
Orbitally-resolved spectral density for the Jeff=1/2

Total spectral function



Calculated spectral density

Jeff=3/2 spectral function

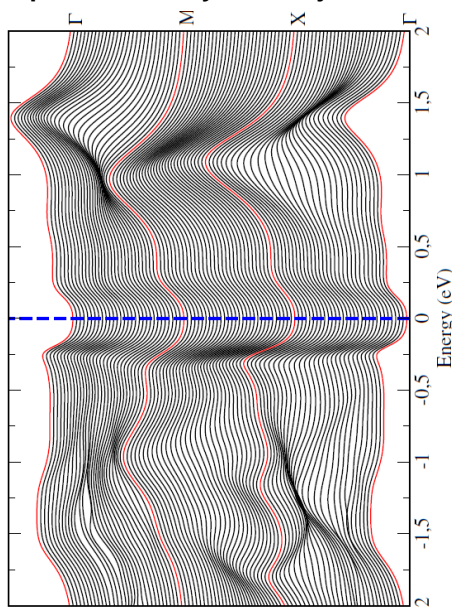


(b) Spectral density for Jeff=3/2

# Spectral function of $\text{Sr}_2\text{IrO}_4$ at 300 K: LDA+DMFT vs. Experiment (Perfetti et al.)

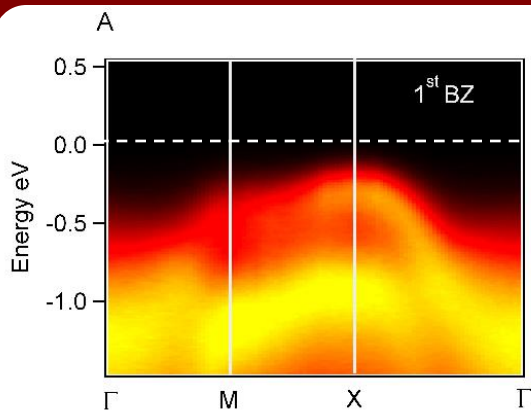
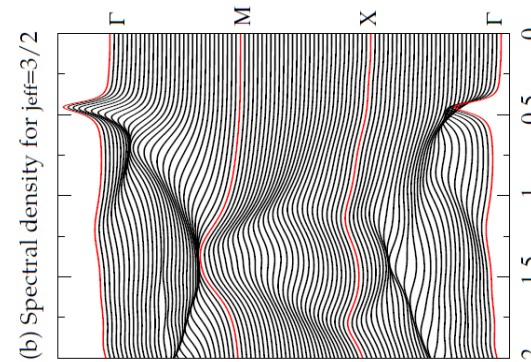
Martins et al., Phys. Rev. Lett. 107, 266404 (2011)

Orbitally-resolved  
spectral density for the  $j_{\text{eff}}=1/2$

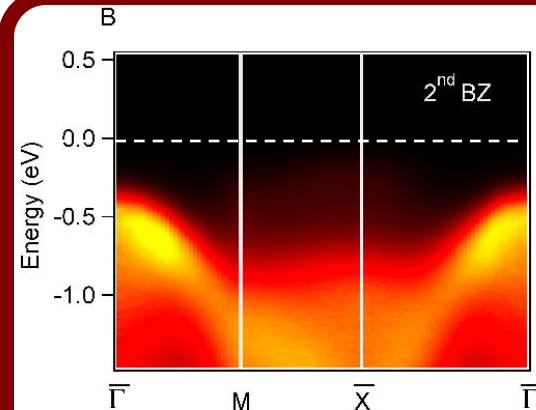
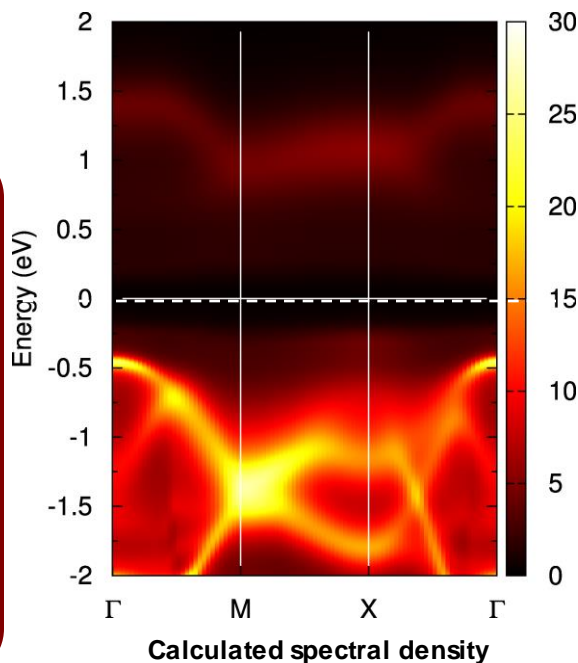


While the  $j_{\text{eff}}=3/2$  orbitals are well described,  
the  $k$ -dispersion of the  $j_{\text{eff}}=1/2$  is missing.

Orbitally-resolved  
spectral density for the  $j_{\text{eff}}=3/2$



Photoemission spectrum at 300 K  
In the 1st Brillouin Zone

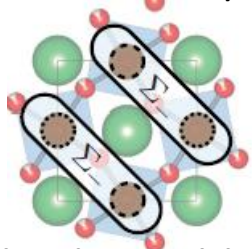


Photoemission spectrum at 300 K  
In the 2nd Brillouin Zone

# Spectral function of $\text{Sr}_2\text{IrO}_4$ at 300 K: DFT + cluster-DMFT vs. Experiment

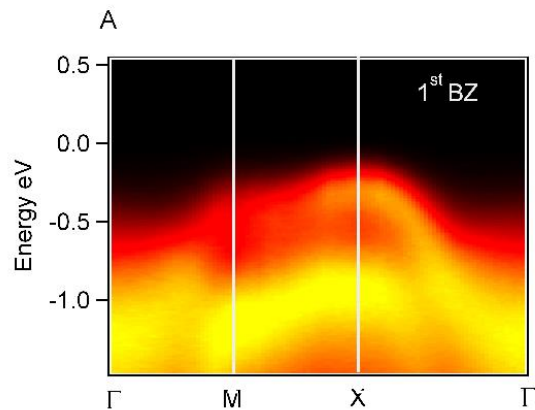
Martins et al., *Phys. Rev. Lett.* 107, 266404 (2011)

Cluster-DMFT  
(2-site cluster)

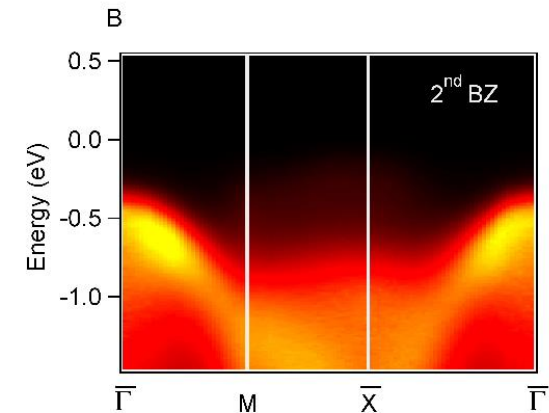
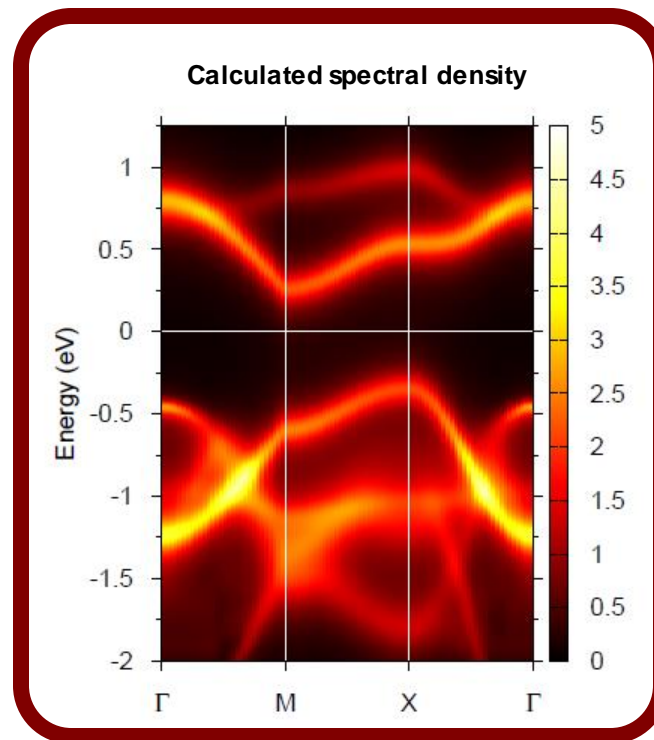


solved in Hubbard  $i$

By taking into account inter-site (non-local) correlations,  
excellent agreement with experiments



Photoemission spectrum at 300 K  
In the 1st Brillouin Zone



Photoemission spectrum at 300 K  
In the 2nd Brillouin Zone

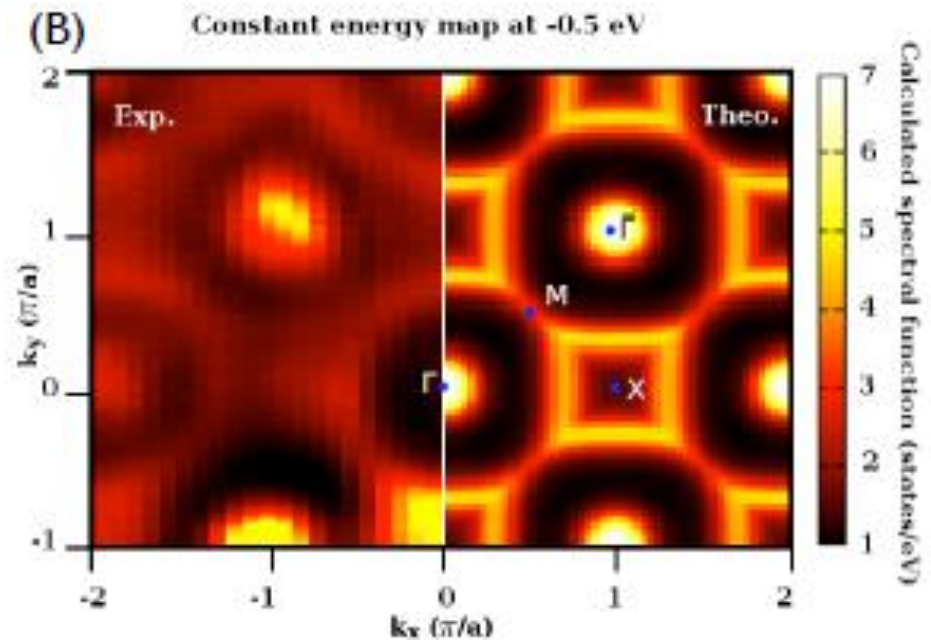
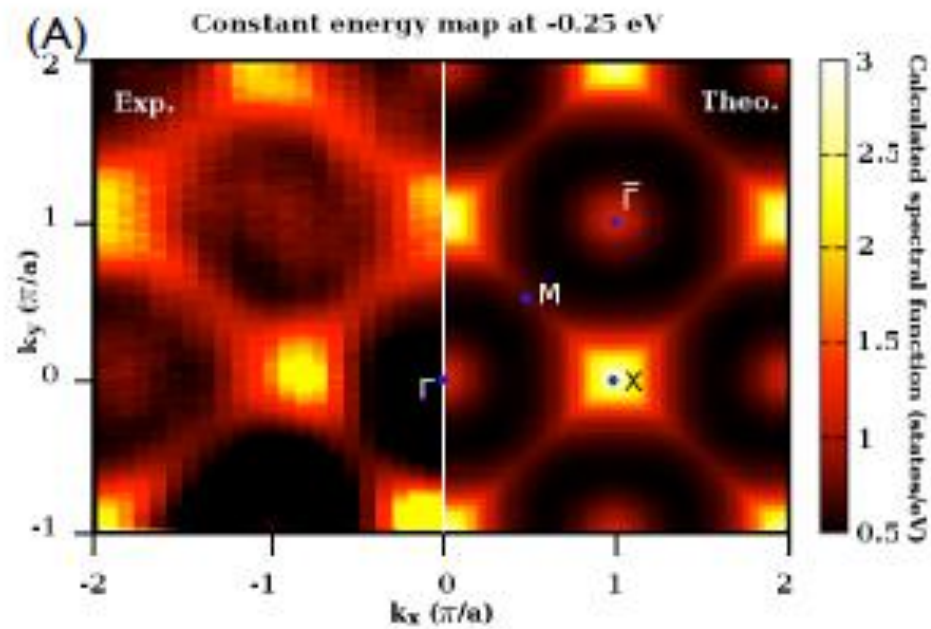
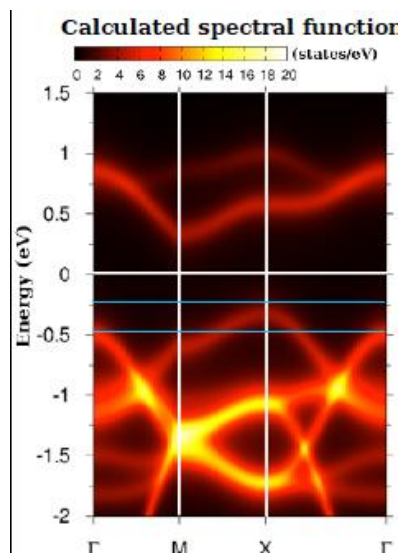
ARPES by L. Perfetti et al.

Theory: Keep 3/2 states from DFT+DMFT; Cluster-Hubbard I for 1/2 state

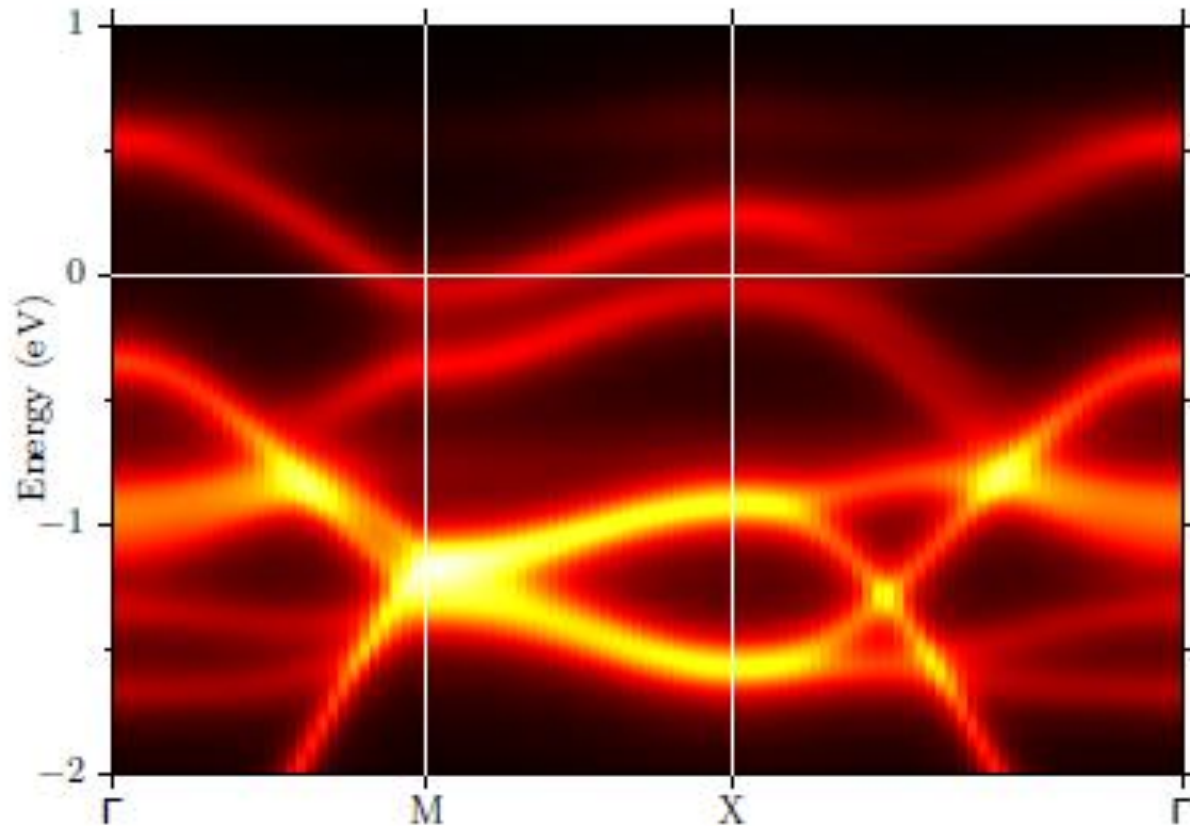
C. Martins, B. Lenz, V. Brouet, F. Bertran, L. Perfetti, S.B. PRMat. Rapid 2018<sup>28</sup>



# Sr<sub>2</sub>IrO<sub>4</sub> -- Constant energy maps: experiment vs. theory



# Doped case: Sr<sub>2</sub>IrO<sub>4</sub> + 10% electron doping



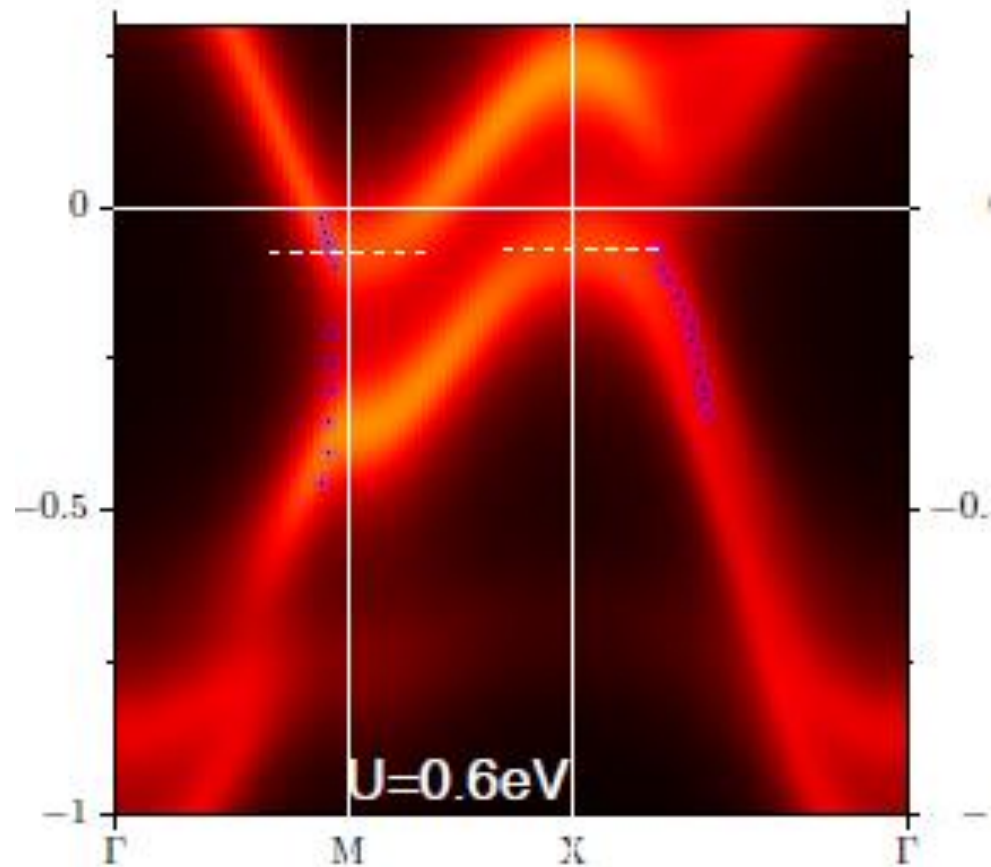
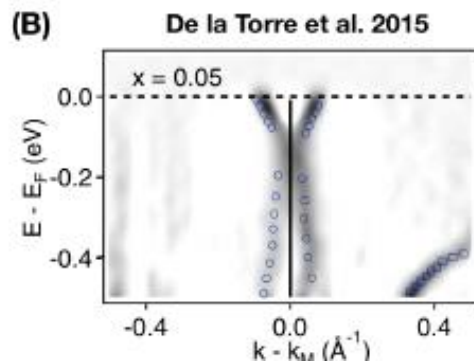
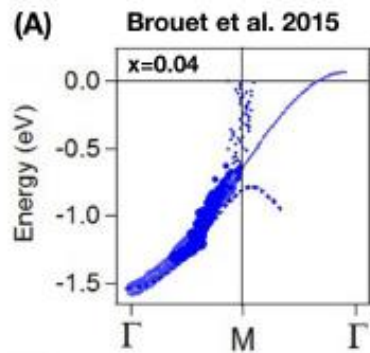
However:  $U$  is reduced by metallic screening!

In practice:  
“VCA”-like solver (no bath sites)

FIG. 3. Calculated momentum-resolved spectral function  $A(\mathbf{k}, \omega + i\eta)$  of 10% electron-doped Sr<sub>2</sub>IrO<sub>4</sub> along the  $\Gamma - M - X - \Gamma$  path. C. Martins, B. Lenz, V. Brouet, F. Bertran, L. Perfetti, S.B. PRMat. Rapid 2018

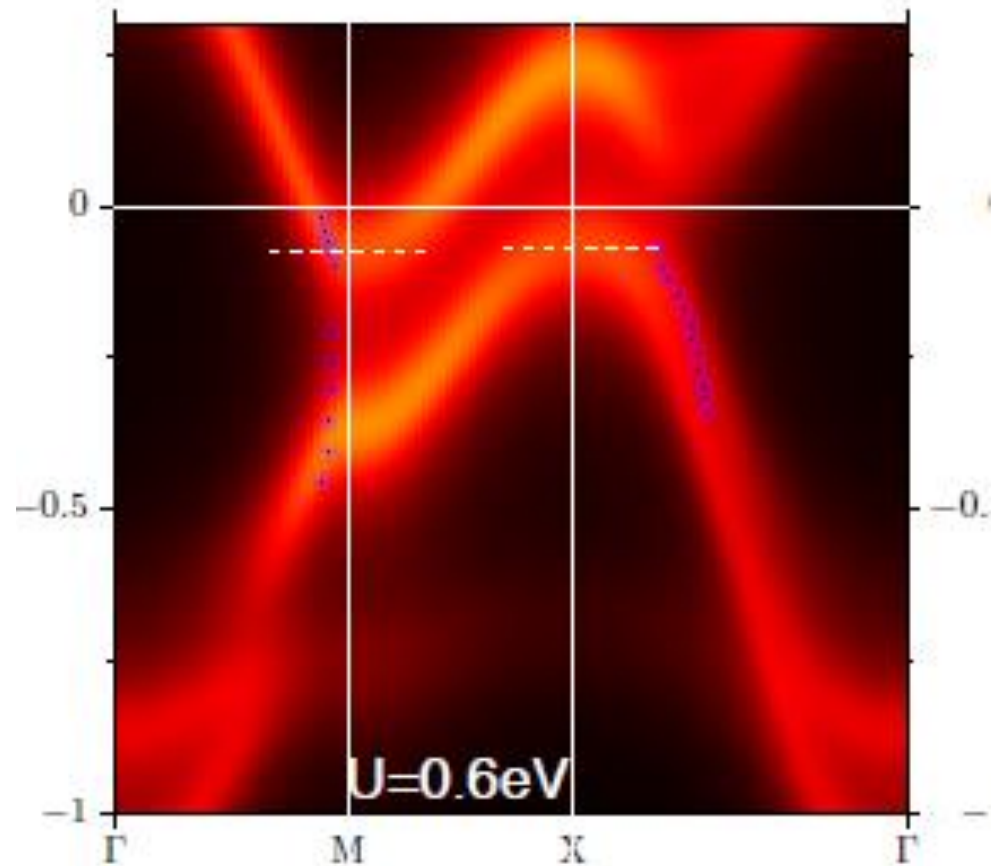
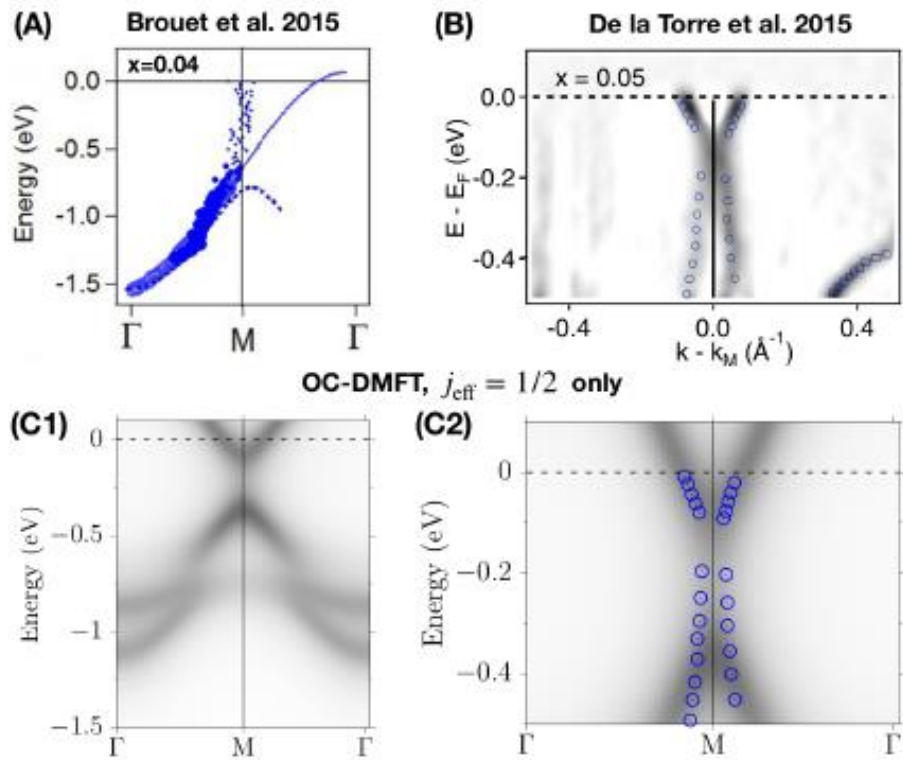


# Doped case: Comparison to ARPES



Blue-dashed: ARPES from  
De la Torre et al. 2015

# Doped case: Comparison to ARPES



Blue-dashed: ARPES from  
De la Torre et al. 2015

# Doped case: Sr<sub>2</sub>IrO<sub>4</sub> + 10% electron doping

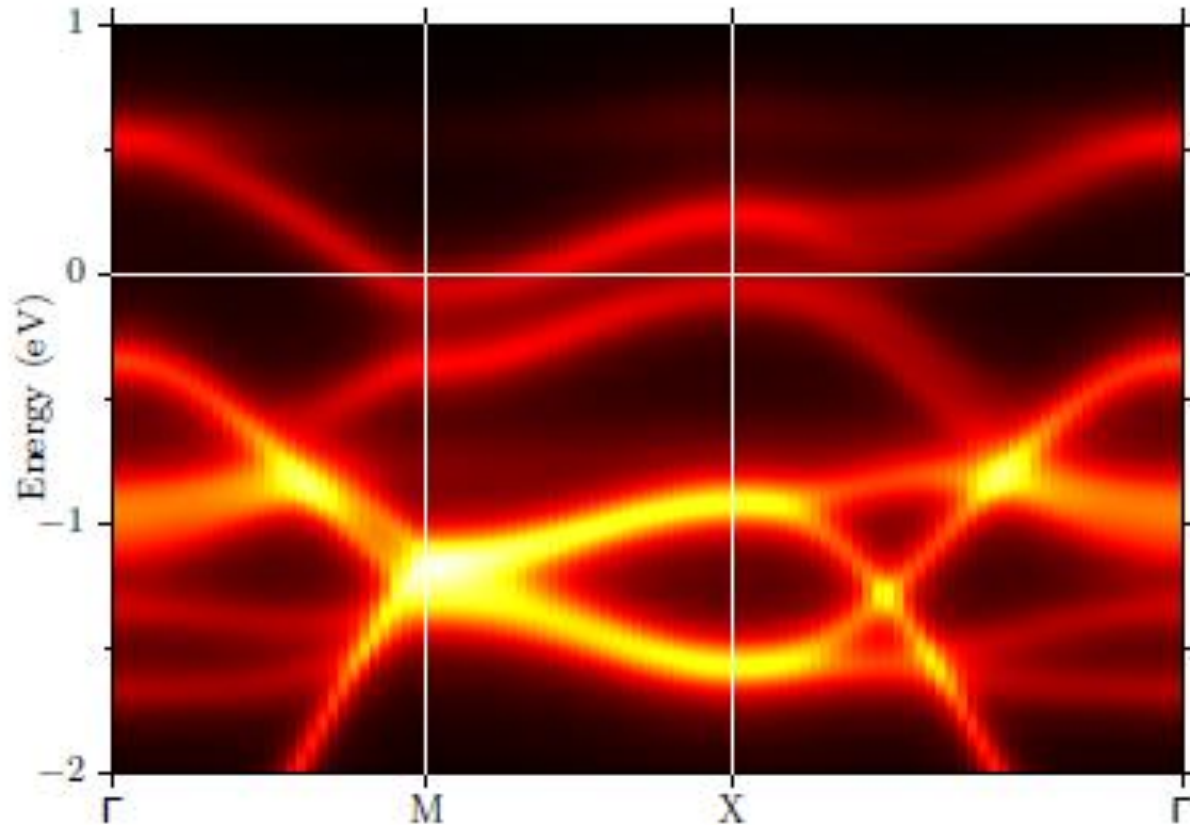
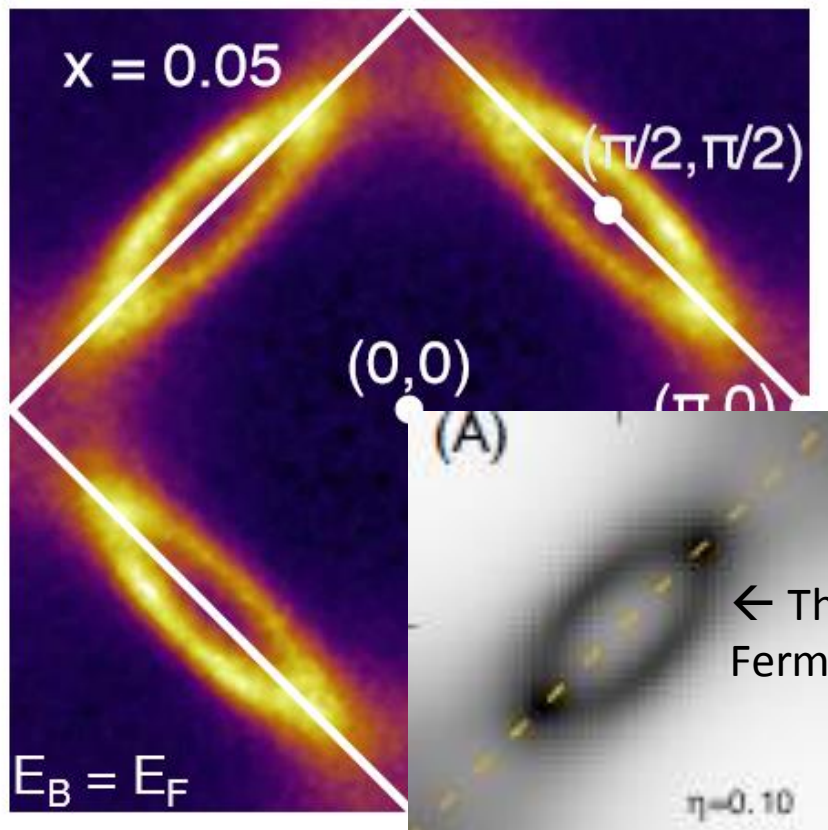


FIG. 3. Calculated momentum-resolved spectral function  $A(\mathbf{k}, \omega + i\eta)$  of 10% electron-doped Sr<sub>2</sub>IrO<sub>4</sub> along the  $\Gamma - M - X - \Gamma$  path. C. Martins, B. Lenz, V. Brouet, F. Bertran, L. Perfetti, S.B. PRMat. Rapid 2018

# Fermi surface of (Sr,La)<sub>2</sub>IrO<sub>4</sub>:



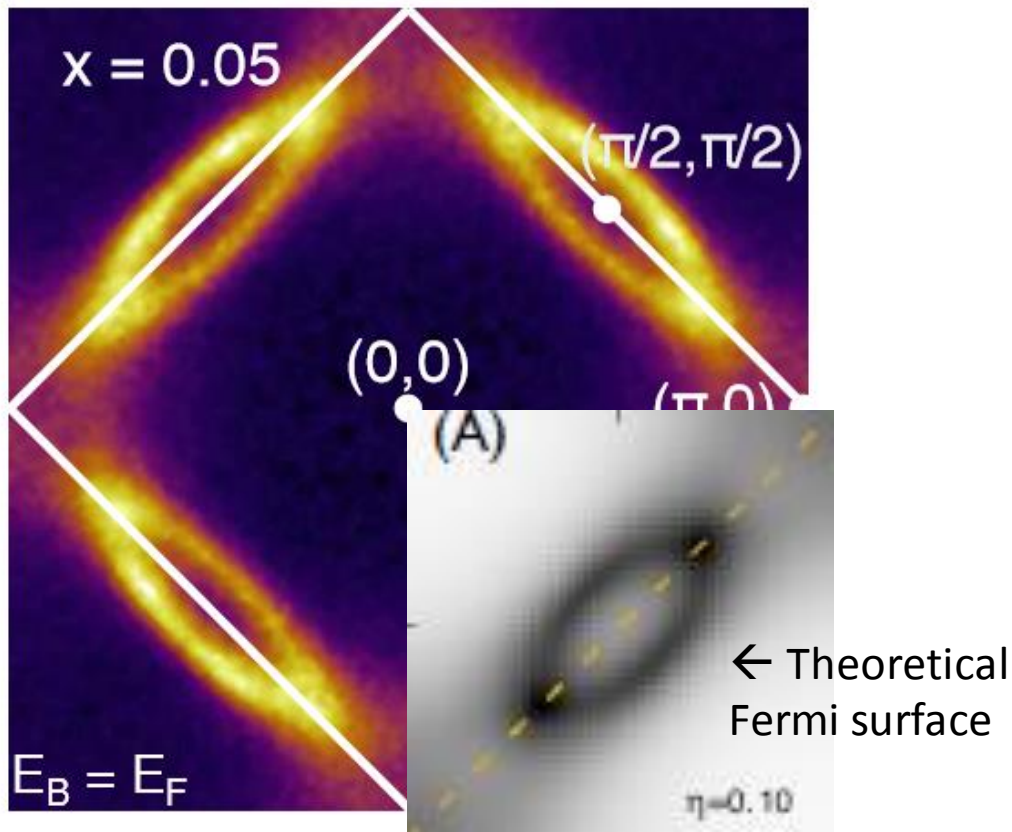
← Theoretical  
Fermi surface

(Martins, Lenz, Perfetti, Brouet,  
Biermann, Phys. Rev. Mat. Rapid (2018))

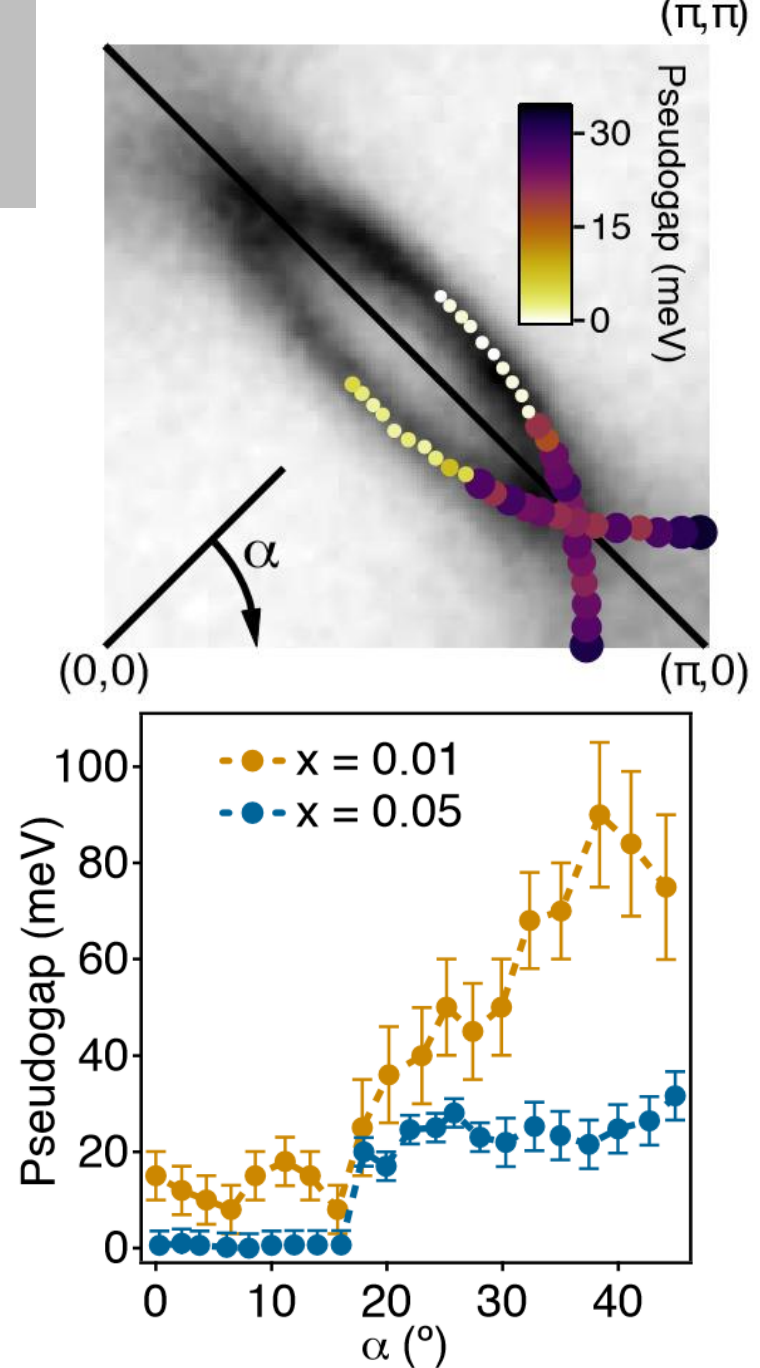
De la Torre *et al.* PRL 2015

See also: Moutenet et al for a QMC  
study of the FS within a model for Sr<sub>2</sub>IrO<sub>4</sub>

# Fermi surface of (Sr,La)<sub>2</sub>IrO<sub>4</sub>:



De la Torre *et al.* PRL 2015

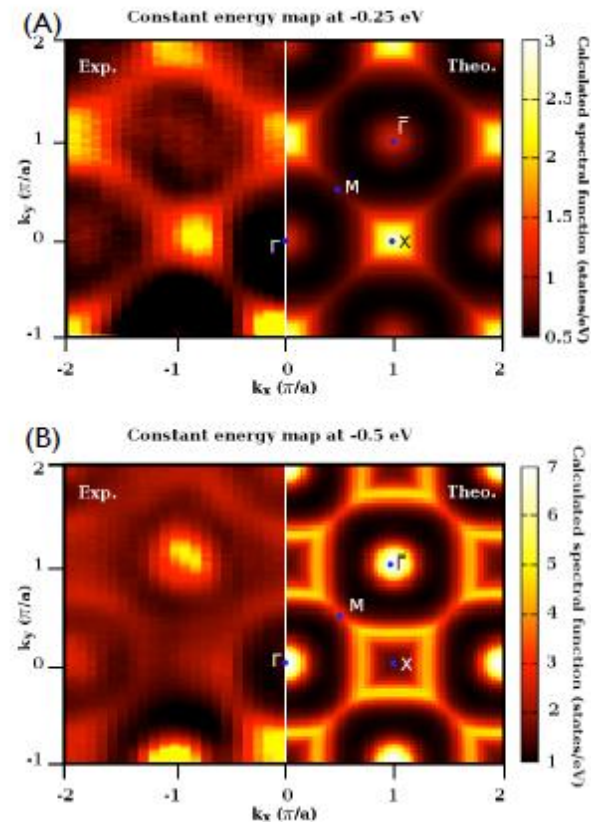




# Summary

”First principles” calculations of spectral properties of iridates: Sr<sub>2</sub>IrO<sub>4</sub>

- Reduced effective degeneracy
- Role of non-local correlations: antiferromagnetic fluctuations crucial (note: not a statement on range!)
- Doped case: good agreement with ARPES!
- Pseudogap-like features a result of several effects: antiferro fluctuations, filled band coming close ...
- “OC-DMFT” a promising cluster scheme



Martins et al., PRL 2011,  
JPCM 2016, PRM (R) 2018



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DE PARIS



## Collaborators and Funding

Sr<sub>2</sub>IrO<sub>4</sub>/Sr<sub>2</sub>RhO<sub>4</sub>  
(PRL 2011 and JPCM 2017)

**C. Martins**, M. Aichhorn, L. Vaugier, S. Biermann

Sr<sub>2</sub>IrO<sub>4</sub> vs. Ba<sub>2</sub>IrO<sub>4</sub>

**C. Martins**, B. Lenz, L. Perfetti, S. Moser

Non-local effects in Sr<sub>2</sub>IrO<sub>4</sub>  
(Phys. Rev. Mat. Rapid 2018)

**B. Lenz**, **C. Martins**, L. Perfetti, V. Brouet

Many useful discussions: Y. Sidis, P. Bourges, J. Jeong

**Thank you for your  
attention !**