Ionization processes of atoms and molecules: a Sturmian approach

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GDR NBODY, Lille, 10 janvier 2020

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Hydrogen spectrum



"challenging the continuum" ... is back !

PLAN

➢ Ionization processes of atoms/molecules → N-BODY

Generalized Sturmian Function (GSF) method implemented to study collision processes

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Ionization of atoms and molecules by photon or particle impact



WHY DO WE WANT DIFFERENTIAL CROSS SECTIONS?

> needed as database for applications (e.g. plasmas or radiobiology)

- test theoretical models in more details
- the collision dynamics the wave functions
- information/test on electronic correlation

SINGLE IONIZATION : (e,2e) on H



Detection in coincidence: Kinematically complete



Pure 3-body problem in final channel

- Solved numerically at the end of century (24 Dec 1999 – Rescigno et al)
- > Agreement between theories and experiments





(double continuum)





FOCUS: Continuum states **NOT** bound states

DOUBLE IONIZATION : (γ,2e)



$$\gamma$$
 + He \longrightarrow He⁺⁺ + e⁻ + e⁻

Detection in coincidence: Kinematically complete FDCS



Double photoionization of He $\gamma + \text{He} \rightarrow \text{He}^{++} + e^{-} + e^{-}$

ECS (McCurdy et al, PRA, 2004)



kdkf



FIG. 3. Real part of direct contribution to the wave function at 20 eV. The panels from top to bottom show the contributions from the kskp, kpkd, and kdkf partial waves.

GSF (Randazzo et al, EPJD, 2015)



Double photoionization (γ,2e) of He

TDCS for photon energy 20 eV above threshold

Equal energy sharing $E_1 = E_2 = 10 \text{ eV}$; different θ_1 values

(Randazzo et al, EPJD, 2015)



Figure 8: (Color online) TDCS for 20 eV and equal energy sharing at various geometries. Circles: absolute experimental values by Brauning *et al.* [14]. <u>Red dashed line: ECS</u> calculation [22]. Black solid line: present GSF results.

Absolute scale

Save 75% of memory storage



Complete experiments \rightarrow multiply differential cross sections \rightarrow transition matrix element $\langle \Psi_i | V | \Psi_f \rangle$ 6D, 9D integrals

Need description (wave function) of charged particles in Coulomb interaction

• BOUND STATES (target)

• CONTINUUM STATES (single or double) (e.g. two electrons escaping from a positive nucleus – THREE-BODY)

N-Body \rightarrow few body (2, 3 or 4) Solve NR time-independent Schrödinger equation

Continuum states



In the final state the ejected electron will « see » asymptotically a charge 1

Central potential U(r) with a Coulomb tail : -Z/r

$$R_l \to e^{ikr - i\frac{Z}{k}\ln(2kr)}$$



Contrary to the bound case here **the energy E is known**. Highly oscillating and long-range. How to represent the continuum? The choice is crucial for efficiency !

GENERALIZED STURMIAN APPROACH

LIKE most other numerical approaches it uses two-body basis functions

BUT

UNLIKE other methods the basis functions have adequate asymptotic behavior

<u>Note:</u> three body \rightarrow also in hyperspherical coordinates

$$ho = \sqrt{r_1^2 + r_2^2} \qquad lpha = \arctan(r_1/r_2)$$

Generalized Sturmian Functions (two-body GSF): S_{n,I}(r)

- > are (numerical) solutions of a Sturm-Liouville differential equation
- ➢ form a complete and discrete set → BASIS SET (index n) (spectral method)
- have a unique and appropriate asymptotic behavior

(with correct energy for continuum states or expected decay for bound states)



$$S_{n,l}(r) \rightarrow e^{ikr - i\frac{Z}{k}\ln(2kr)}$$

- > can be constructed to properly represent the electron-nucleus cusp
- concentrate the effort in the inner part where interaction takes place 14

BOUND (E<0)

Same boundary conditions (for all n)

$$S_{n,l}(r) \rightarrow e^{-ar}$$



Left: Energy eigenfunctions: widely spread. Right: Potential eigenfunctions. The expansion power can be focused in the area of interest.

Bound states

GROUND STATE

He: BEST ground state energy with uncorrelated product

-2.903 712 820 -2.903 724 377 [14]

Exotic systems e.g. positronium ion

Table 7.7 Partial-wave convergence of the ground-state energy of Ps^{-} system. We used 35 radial functions per coordinate and for each partial-wave I_i

Ps ⁻ Ground-state energy				
li	E _{li}			
0	-0.257 240 143			
1	-0.260 105 390			
2	-0.261 496 276			
12	$-0.262\ 002\ 458$			
Exact ⁶¹	-0.262 005 070			

EXCITED STATES

Table 7.3 Energy of the first three excited states for singlet states of He for different L with a total of 168 basis functions

n	L	Present work	Ref. 57
3	2	-2.0556110426	-2.0556207328522456
4	3	-2.03125512987	-2.0312551443817490
5	4	-2.020000709670	-2.0200007108985847
6	5	-2.0138890317669	-2.0138890347542797

DOUBLY EXCITED STATES

Table 7.4 Energies for the Rydberg series of the singlet S states of He, using 30 GSF per electron

(N, k) _n	GSF		Ref. 58	
	위(E)	3 (E)	ભ(E)	ે(E)
$(2,1)_2$	-0.777876955	-0.002060106	-0.777867636	-0.002270653
$(2, -1)_2$	-0.621817695	-0.000106535	-0.621927254	-0.000107818
$(3,2)_3$	-0.351827523	-0.001406250	-0.353538536	-0.001504906



Left: Energy eigenfunctions, continuous spectrum. Right: Potential eigenfunctions. When outgoing (incoming) conditions are chosen, they constitute a discrete set, even when energies are positive. The potential $\mathcal{V}_{q}(r)$ regulates the spatial extension of the basis expansion capabilities.



GENERALIZED STURMIAN APPROACH

- Efficient basis (smaller computational resources)
 Adequate asymptotic conditions already built-in

 concentrate the effort in the inner part where interaction takes place
- > ADVANTAGE: the extraction of amplitudes and differential cross sections directly from the scattering wave function Ψ^+_{sc} (no need to evaluate a transition matrix element $\langle \Psi_i | V | \Psi_f \rangle$) 19

Applications to ionization processes

Single continuum (two-body)

- **1)** Single photoionization of atoms and molecules (γ ,e)
- 2) Single electron impact ionization of molecules (e,2e)
- 3) Photodetachment of anions (y,e)

4) Two-photon single ionisation (2γ,e)

$$\Psi_{sc}^+ = \sum_n a_n^l \ S_{n,l}^+(r)$$

20

Double continuum (three-body)

- **5)** Double photoionization of He $(\gamma, 2e)$ pure 3-body
- 6) Double electron impact ionization of He (e,3e) pure 4-body
- 7) Double proton impact ionization of He (p,p2e) pure 4-body
- **8) Double photoionization of molecules** (γ,2e)

$$\Psi_{sc}^{+}(\mathbf{r}_{1},\mathbf{r}_{2}) = \sum_{L \ l_{1}l_{2}} \sum_{n_{1}} \sum_{n_{2}} a_{n_{1}n_{2}}^{l_{1}l_{2}L} \mathcal{A} \frac{S_{n_{1}l_{1}}^{+}(r_{1})}{r_{1}} \frac{S_{n_{2}l_{2}}^{+}(r_{2})}{r_{2}} \mathcal{Y}_{l_{1}l_{2}}^{L0}(\widehat{\mathbf{r}}_{1},\widehat{\mathbf{r}}_{2})$$

What about molecules?

- N-BODY
- Multicenter nature
- Average over angular orientation

Calculations on small polyatomic molecules like H_20 , CH_4 , NH_3

Granados et al, Adv Quantum Chem, 2016, **73**, 3 C. Granados and L.U. Ancarani, Eur.Phys.J D, 2017, **71**, 65 Randazzo et al, PRA, 2020



 $H_2O \quad (1a_1)^2(2a_1)^2(1b_2)^2(3a_1)^2(1b_1)^2$ 3 nuclei + 10 electrons : approximations (BO, FC, OCE, ...) → few body problem

1) Single photoionization

 \bigwedge

Μ



M⁺

Photoelectron spectrum



 $\gamma + M \longrightarrow M^+ + e^-$

Inner valence

 $3a_1$ of H₂O

(Granados et al, Adv Quantum Chem, 2016)

2) Single ionization by electron impact



 $e^{-} + M \rightarrow M^{+} + e^{-} + e^{-}$

(e,2e) experiment



Momentum transfer: $\mathbf{K}=\mathbf{q}=\mathbf{k}_{i}-\mathbf{k}_{a}$



 $1b_1 \text{ and } 3a_1$ $E_i=250 \text{ eV}$ $E_b=10 \text{ eV}$ $\theta_a=-15^\circ$

EXPT : Milne-Brownlie et al PRA 69 (2004) 032701

3) Double photoionization

First measurements 2018 (photons of 63 and 65 eV) !! First calculations 2019 (Randazzo et al, PRA, 2020)



Take home message

- Ionization processes
 - → N-body problem involving continuum states of given energy (highly oscillating + long range)
- Generalized Sturmian Functions (GSF)
 - computationally efficient because appropriate boundary conditions are built in the basis elements
 → concentrate the effort in the region where it is interesting
 - implemented to describe the single and double continuum of atoms and molecules

Thank you for your attention !